



Imageries Globale et Régionale et Applications Géodynamiques

Jean-Paul Montagner

Dept. Sismologie, I.P.G., Paris; France

Data

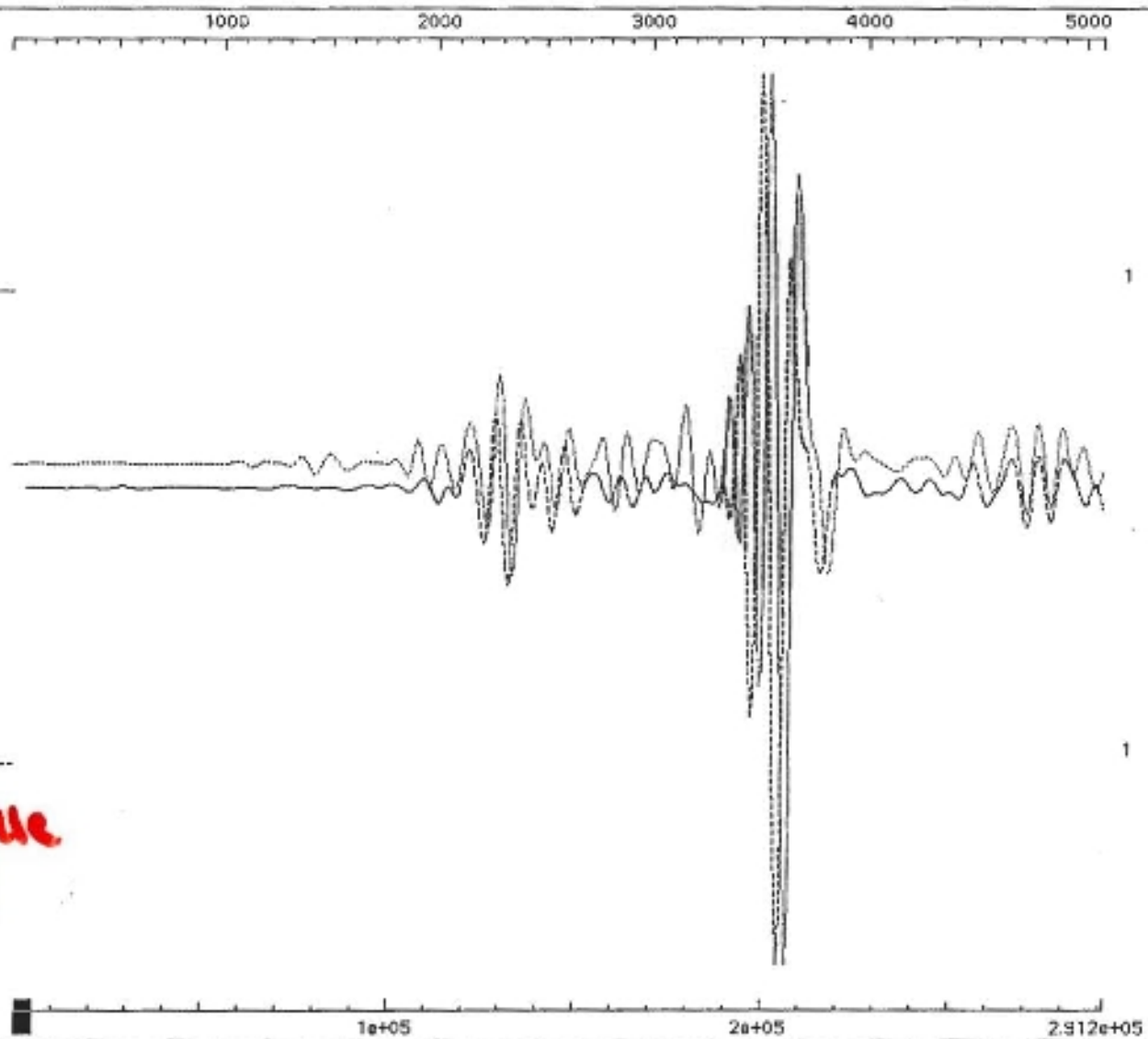
CAN_VHZ19952110511.ah

Donnée

Synthetics

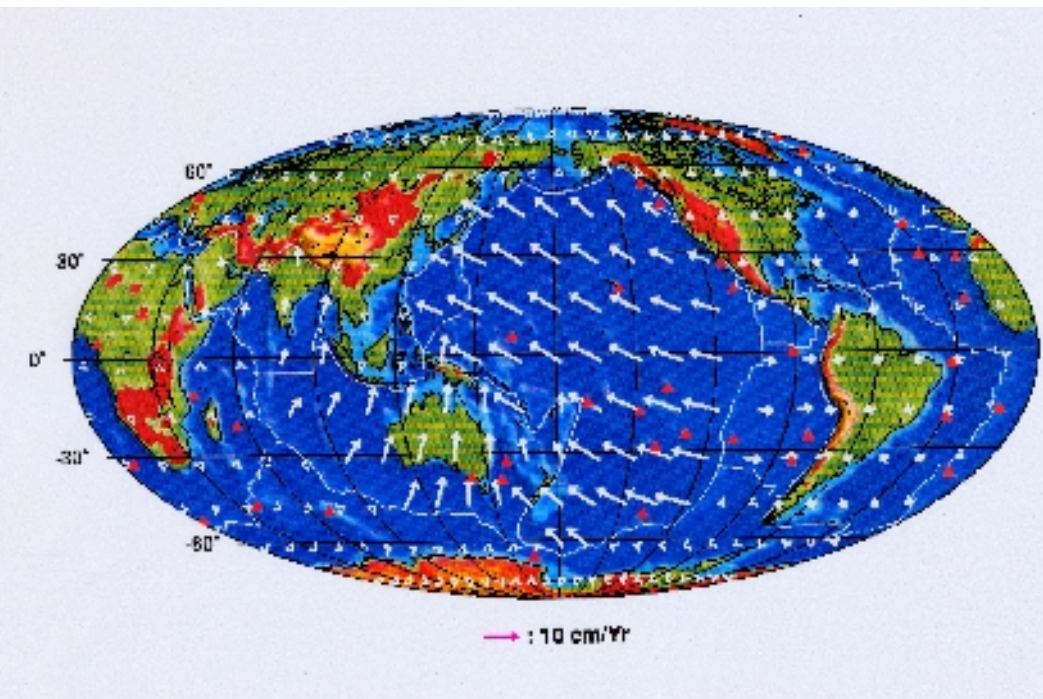
HARMO/cansyr0-6.ah

Synthétique

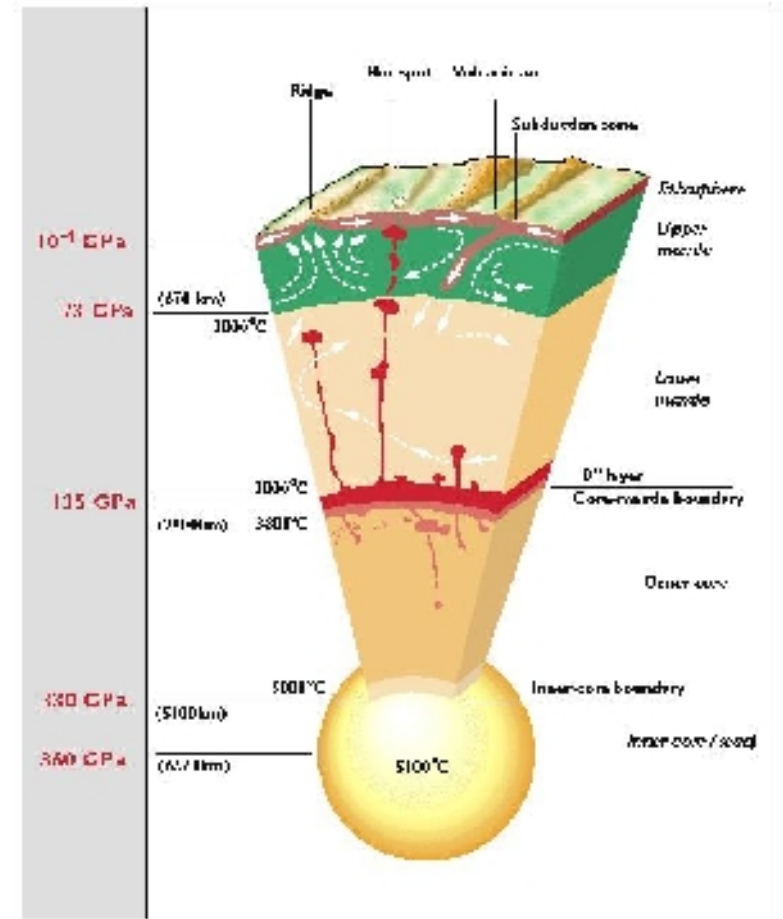


Structure of the Earth

Plate tectonics



Mantle Convection



Tomographic Technique

- **Forward Problem:** Theory $\mathbf{d}=\mathbf{g}(\mathbf{p})$

\mathbf{d} data space, \mathbf{p} parameter space

- Reference Earth model \mathbf{p}_0 :

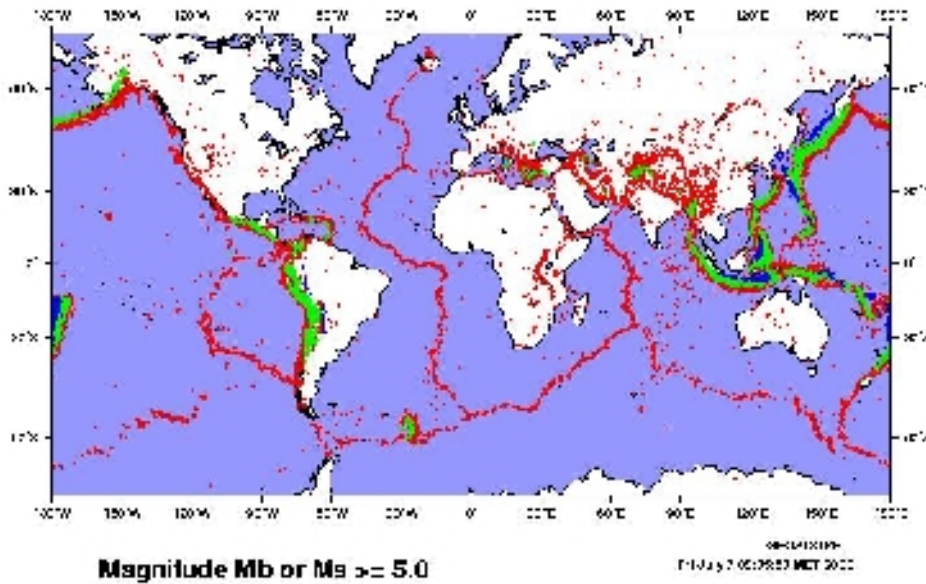
$$\mathbf{d}_0 = \mathbf{g}(\mathbf{p}_0)$$

- Kernels $\partial\mathbf{g}/\partial\mathbf{p}$
- Cd function (or matrix) of covariance of data

- **Inverse Problem:** $\mathbf{p}-\mathbf{p}_0 = \mathbf{g}^{-1} (\mathbf{d}-\mathbf{d}_0)$

- C_{p0} a priori Covariance function of parameters
- C_{pf} a posteriori Covariance function of parameters
- R Resolution

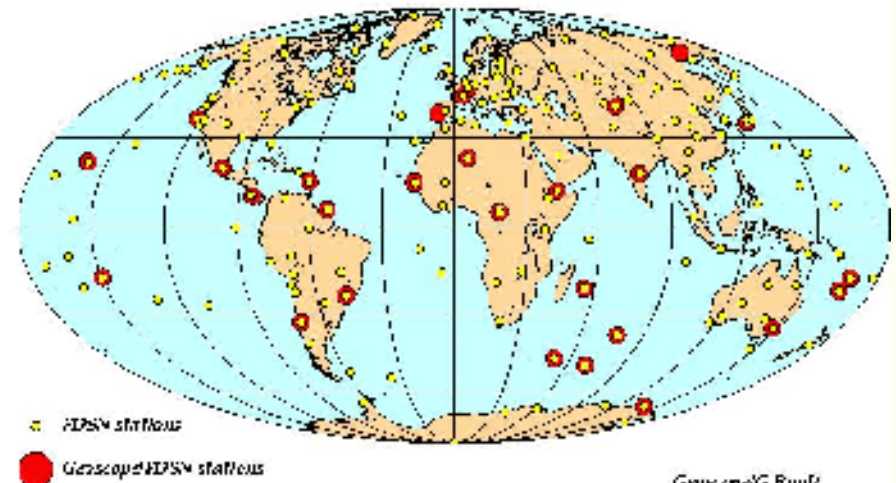
Global seismicity 1928-1999



Des récepteurs

Des sources

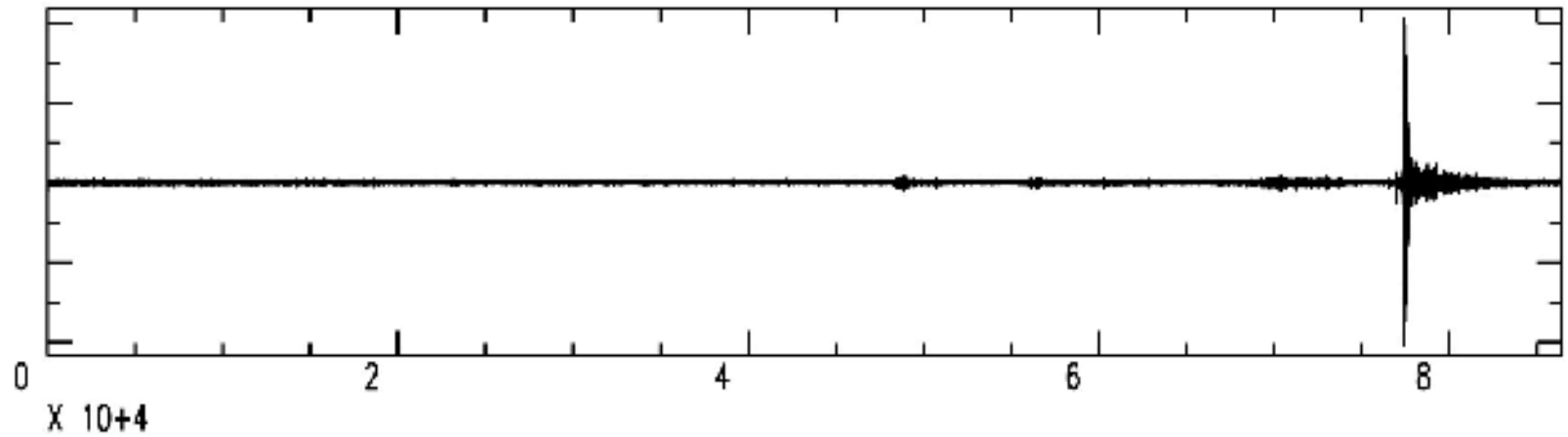
GEOSCOPE stations and FDSN stations



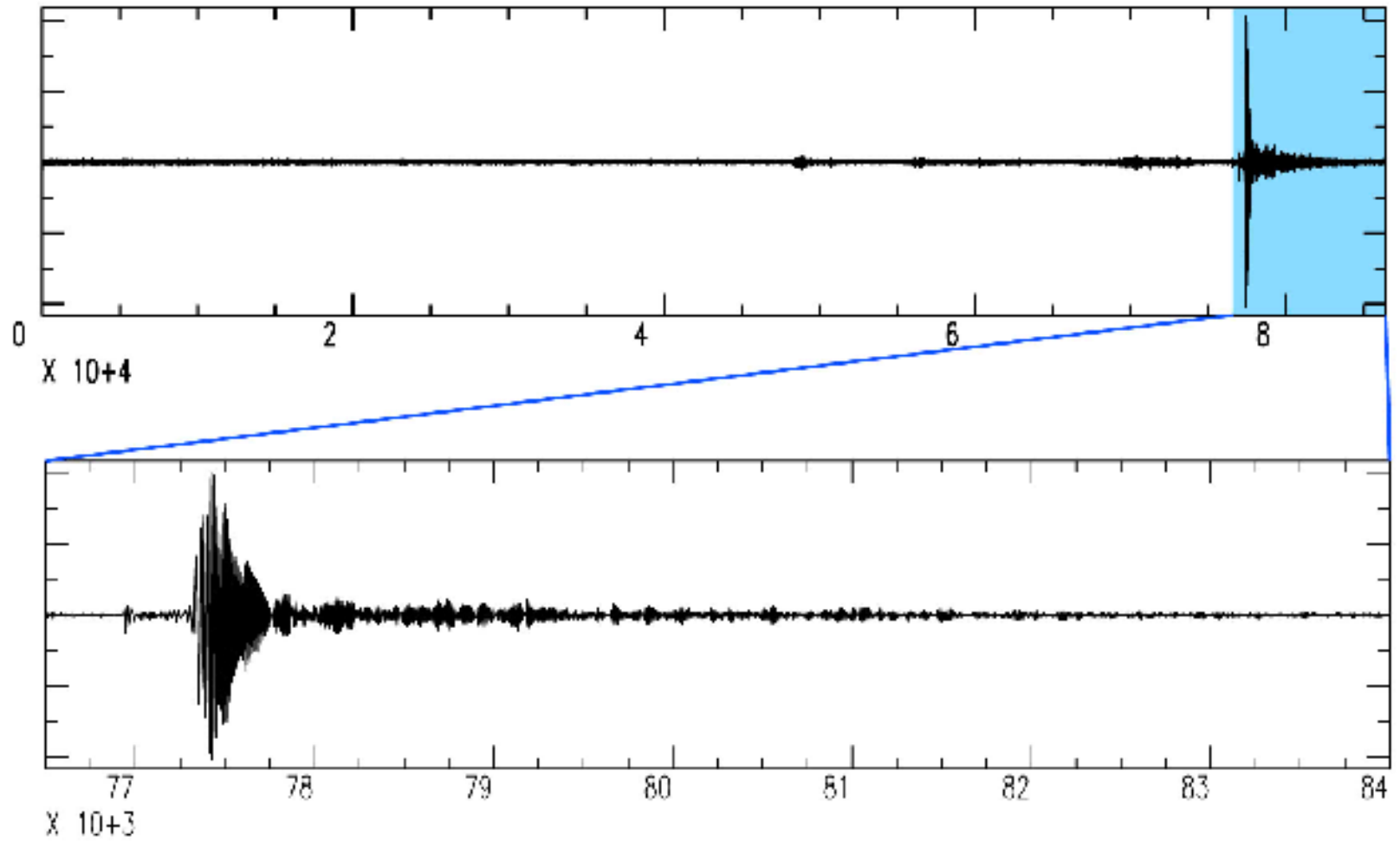
Différents types de données

- Ondes de volume (temps d'arrivée)
P, S,
- Ondes de surface (Rayleigh, Love)
- Modes propres de la terre
- “Bruit de fond” (1-20s microsismique,
>200s “seismic hum”)

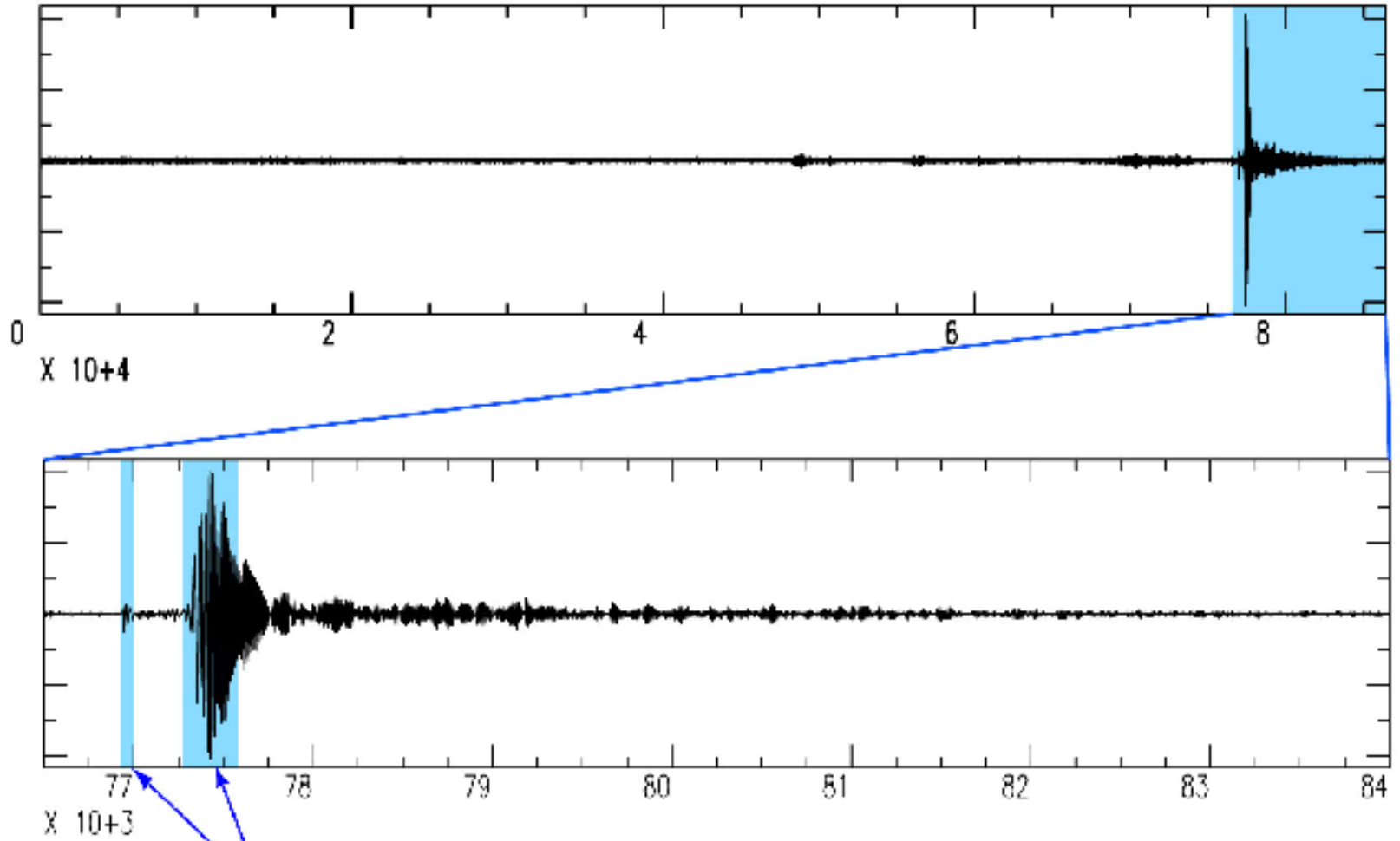
one day of seismic record



one day of seismic record

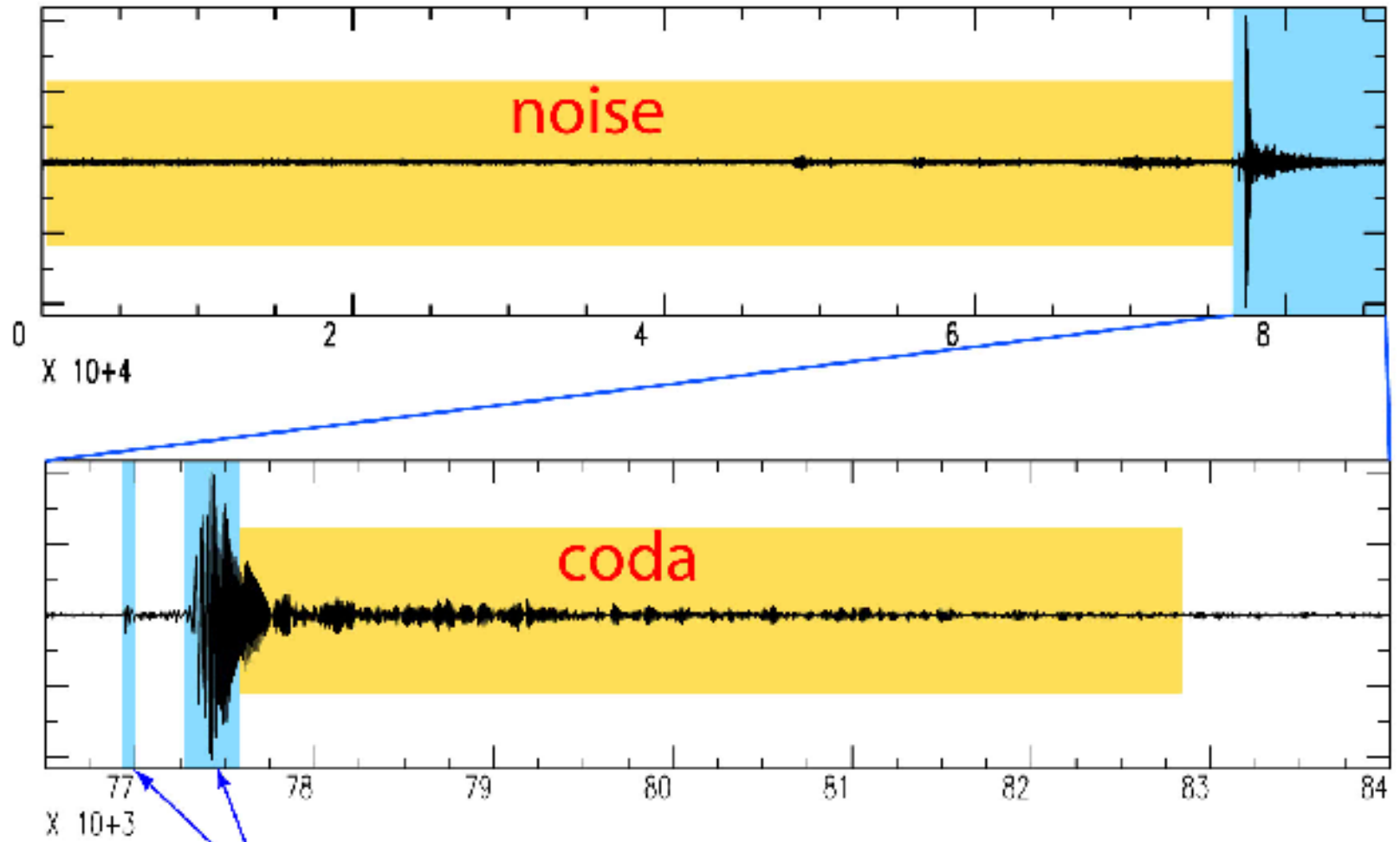


one day of seismic record



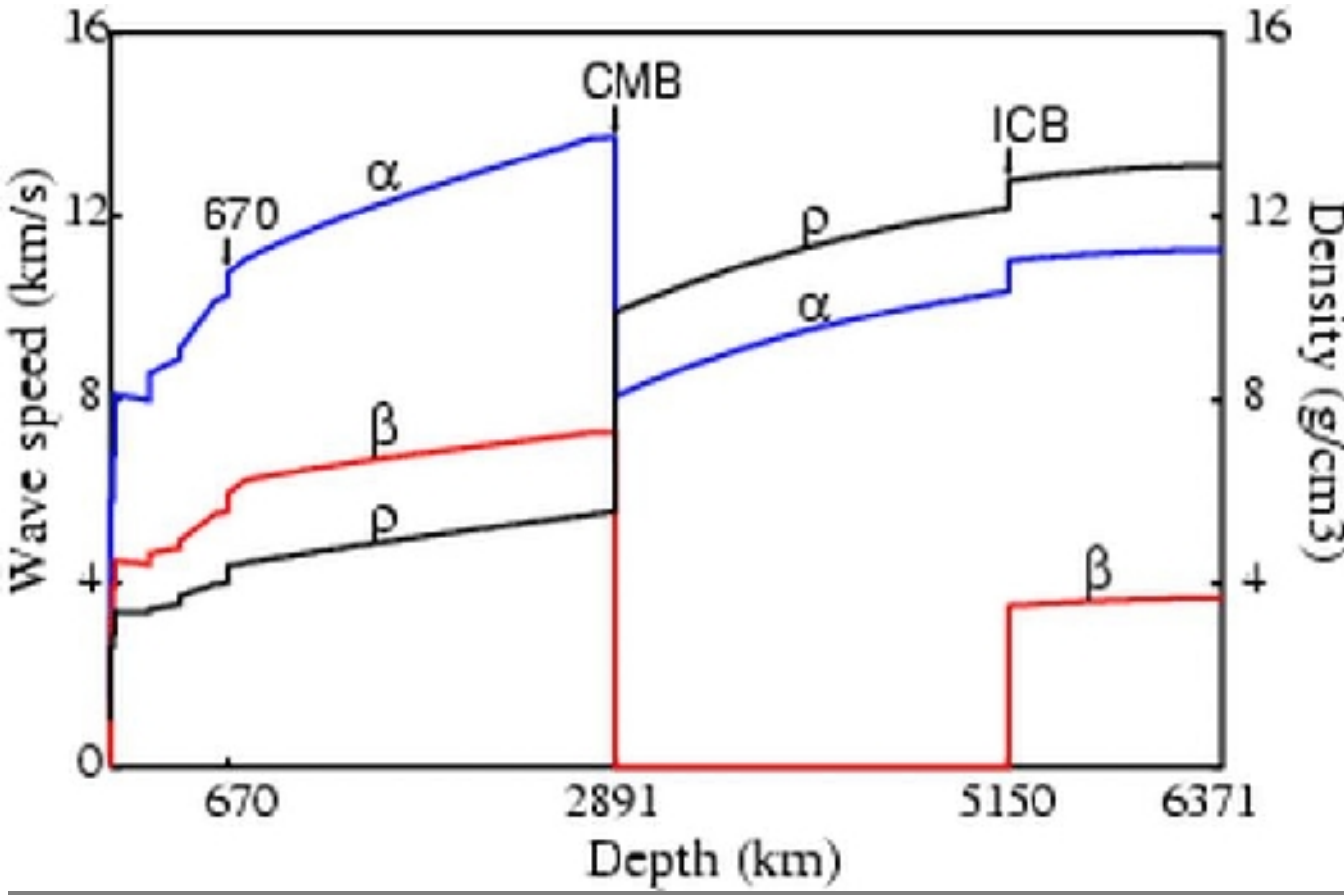
ballistic waves used in traditional tomography

one day of seismic record



ballistic waves used in traditional tomography

1D- Reference Earth Model: PREM

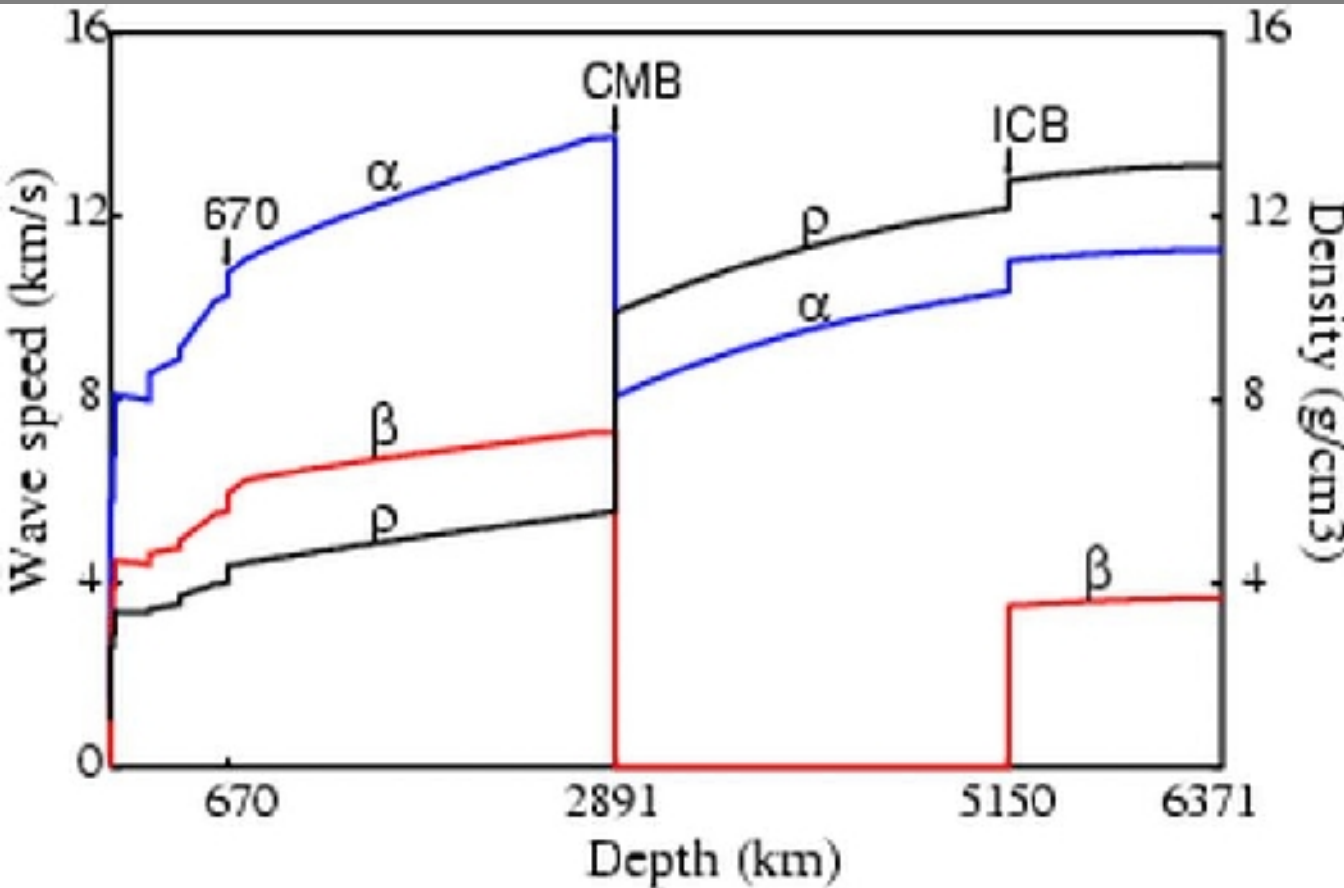


- Normal mode data
- Body wave data (travel times)

$\nabla \rho, V_p, V_s$

$\nabla \rho, V_{ph}, V_{sv}, V_{pv}, V_{sh}, \eta$

1D- Reference Earth Model: PREM

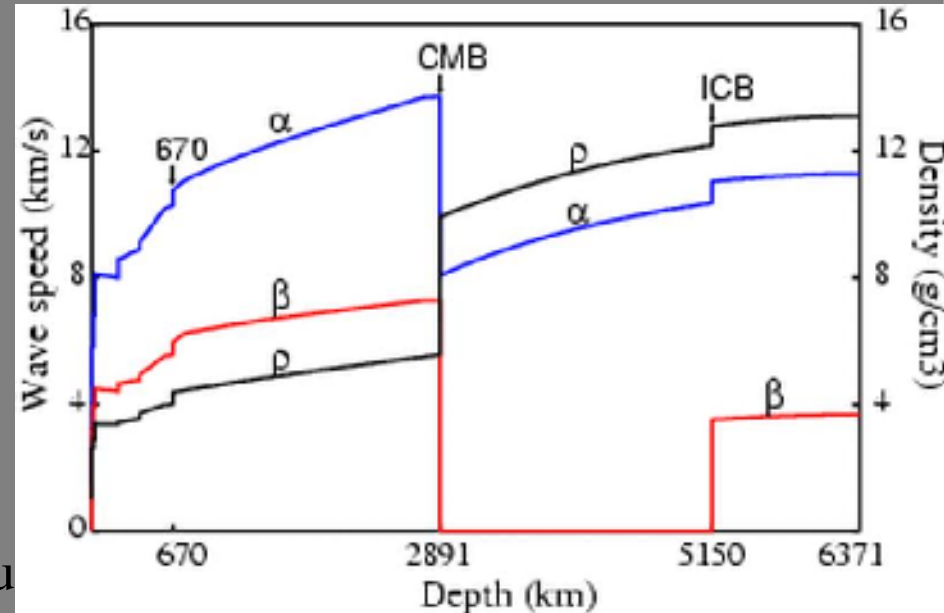


- Crust
- Lithosphere
- Asthenosphere
- Transition zone
- Lower mantle
- D'' layer
- Outer core
- Inner core

1D- Reference Earth Model

- Normal mode calculation
(eigenfrequencies ${}_n\omega_k$
eigenfunctions ${}_n u_l(r,t)$)
- Synthetic Seismograms by
normal mode summation
($k=\{n,l,m\}$).

$\mathbf{u}(\mathbf{r},t)$ Displacement at point \mathbf{r} at time t due
to a force system \mathbf{F} at point source \mathbf{r}_s



$$\mathbf{u}(\mathbf{r},t) = \sum_k \mathbf{u}_k(\mathbf{r}) \cos \omega_k t / \omega_k^2 \exp(-\omega_k t / 2Q) (\mathbf{u}_k \cdot \mathbf{F})_s$$

$$\text{Source Term } (\mathbf{u}_k \cdot \mathbf{F})_s = (\mathbf{M} : \boldsymbol{\varepsilon})_s$$

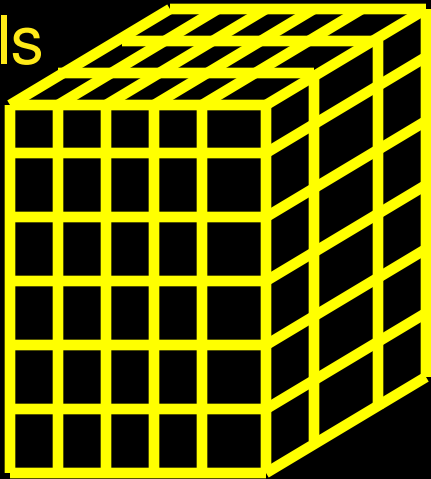
\mathbf{M} Seismic moment tensor, $\boldsymbol{\varepsilon}$ deformation tensor

Parameter Space

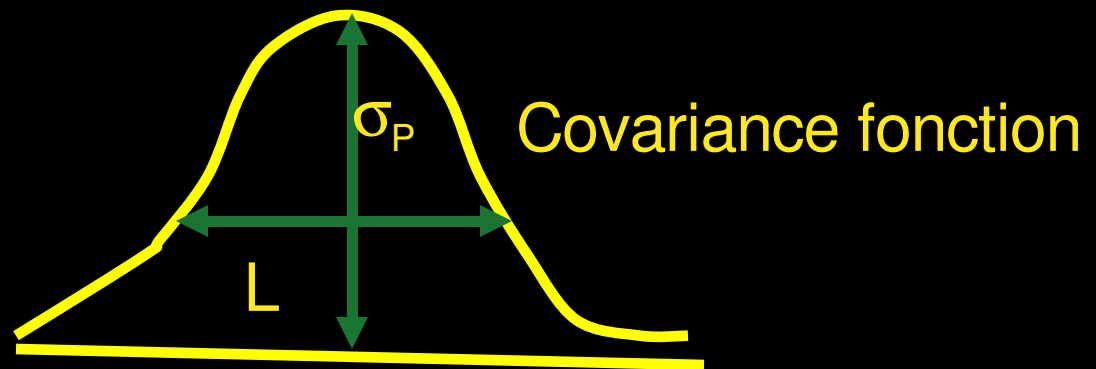
Physical parameters: $\rho + 21$ (13) physical parameters

Geographical parameterization: $\mathbf{p}(r, \theta, \phi)$

Cells



Continuous parameterization



▪ Spherical harmonic expansion

▪ Lateral resolution:

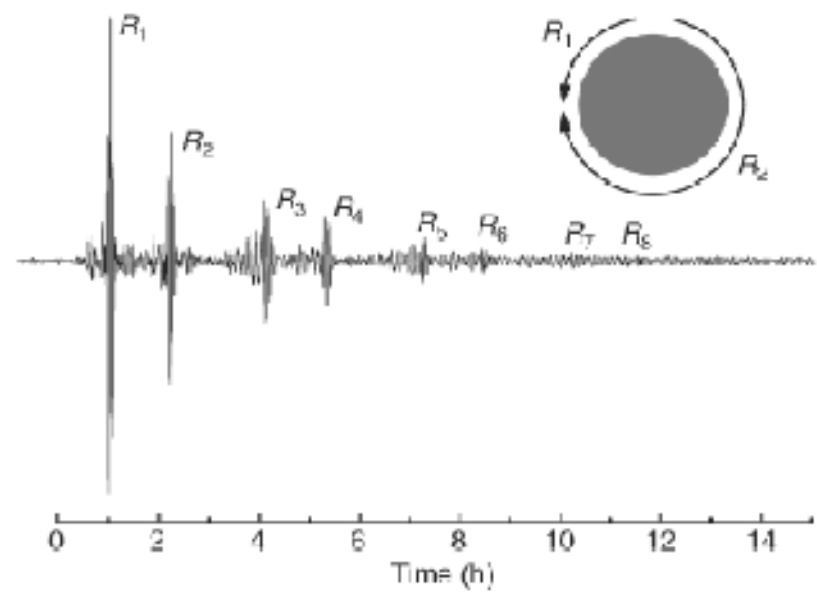
Hor 1000km, Rad 50km $\Rightarrow 500 \cdot 60 \cdot 14 \approx 420,000$ parameters

Différents types de données

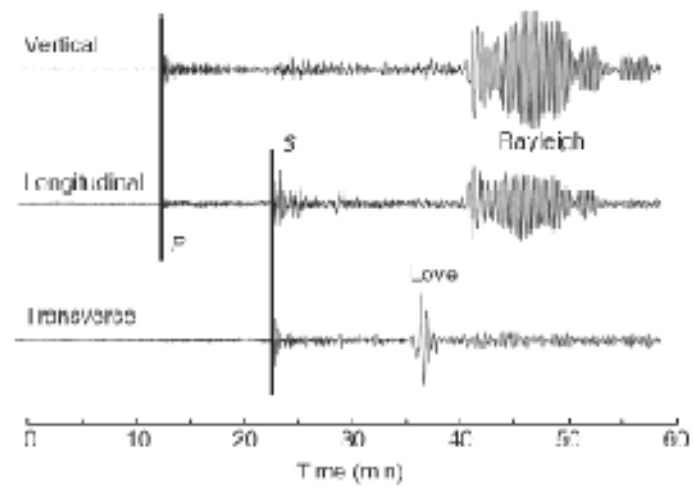
- Ondes de volume (temps d'arrivée)
P, S,
- Ondes de surface (Rayleigh, Love)
- Modes propres de la terre
- “Bruit de fond” (1-20s microsismique,
>200s “seismic hum”)

Data

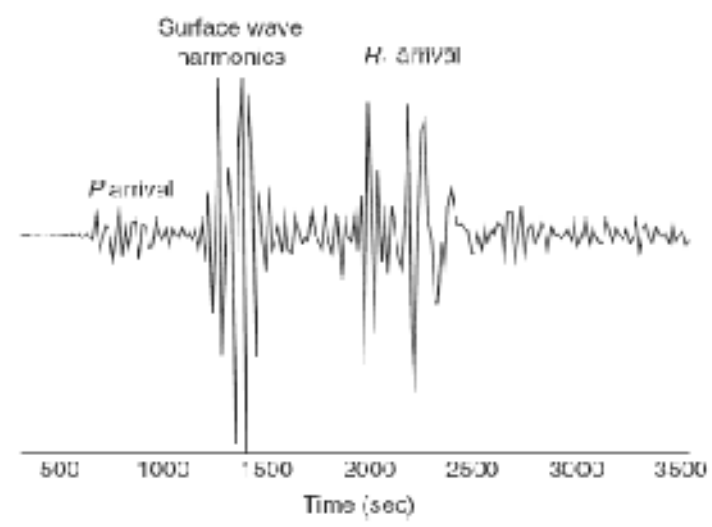
Station: CMB
 Channel: LHZ
 1996/07/11 21:48:39.7 h = 15.0 km $\Delta = 109.7^\circ$ $\phi = 32.0^\circ$
 Burma-China Border Region $M_w = 6.8$



Station: CMB
 1996/07/12 15:13:59.8 h = 15.0 km $\Delta = 83.7^\circ$ $\phi = 7.1^\circ$
 Loyalty Islands Region $M_w = 6.4$



1/23/1997 Southern Colivia
 $M_h = 4$ $M_w = 7.1$ $dep = 276$ km
 GUP Comp L Dist 0415 km

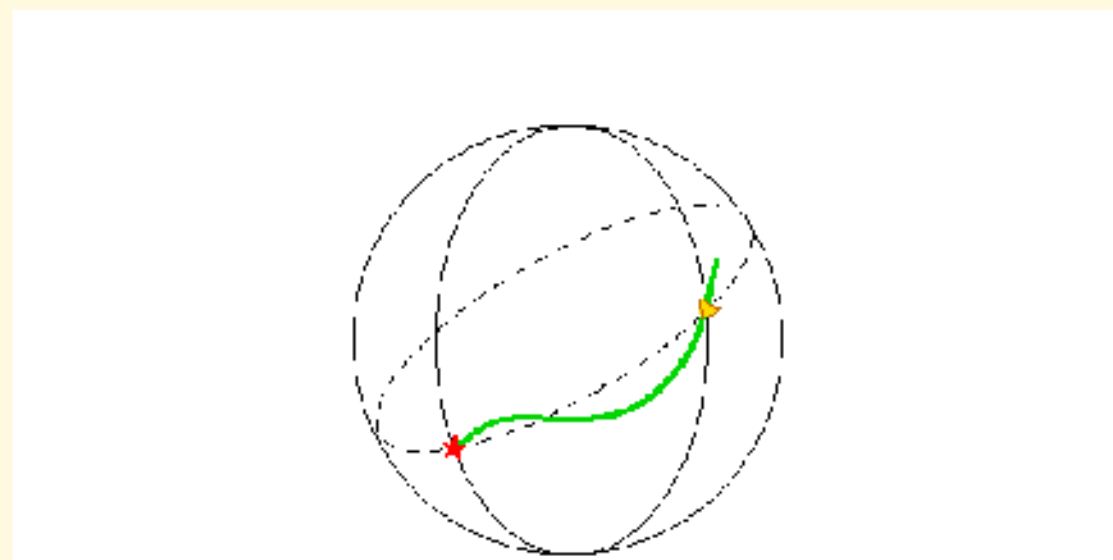


Surface wave full ray theory

Waves locally as plane waves $\lambda \ll \Lambda_\Omega$

$$u(\omega) = \sum_{\text{branches}} \sum_{\text{orbits}} A(\omega) \exp^{-i\phi(\omega)}$$

$$A = A_s A_p A_r \quad \phi = \phi_s + \phi_p + \phi_r$$

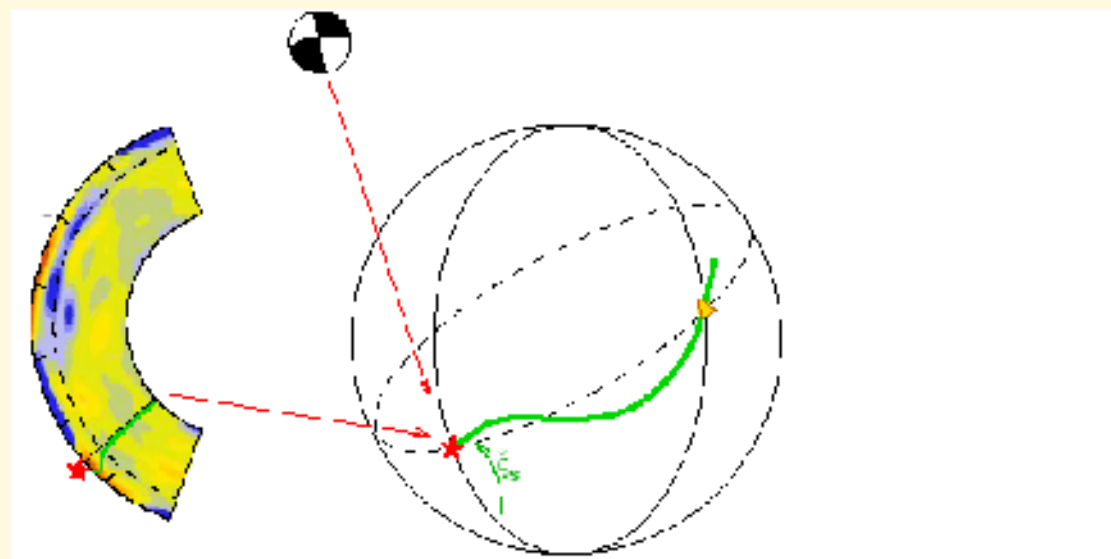


Surface wave full ray theory

Waves locally as plane waves $\lambda \ll \Lambda_{\Omega}$

$$u(\omega) = \sum_{\text{branches}} \sum_{\text{orbits}} A(\omega) \exp^{-i\phi(\omega)}$$

$$A = A_s A_p A_r \quad \phi = \phi_s + \phi_p + \phi_r$$

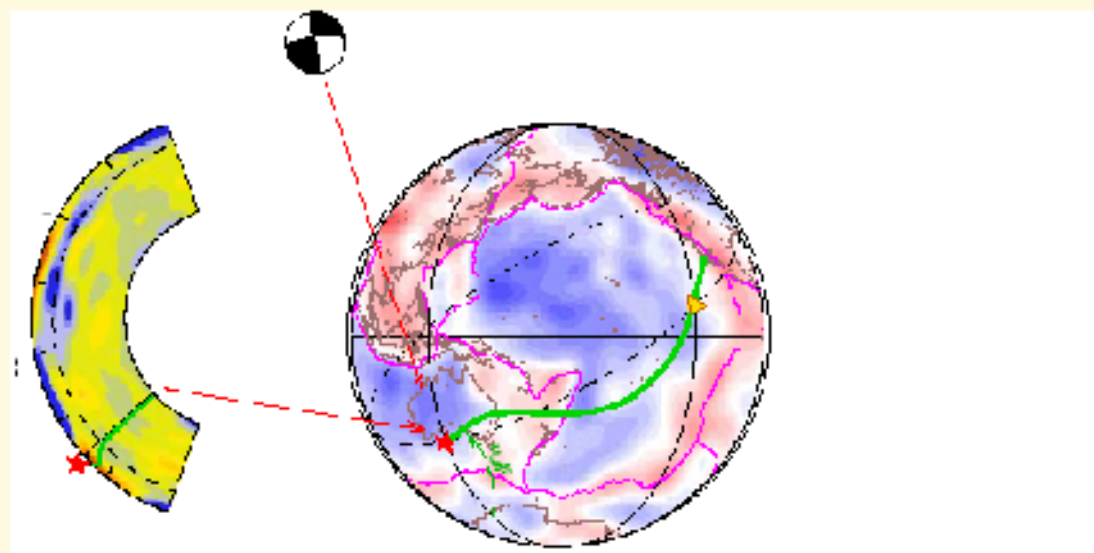


Surface wave full ray theory

Waves locally as plane waves $\lambda \ll \Lambda_\Omega$

$$u(\omega) = \sum_{\text{branches}} \sum_{\text{orbits}} A(\omega) \exp^{-i\phi(\omega)}$$

$$A = A_s A_p A_r \quad \phi = \phi_s + \phi_p + \phi_r$$

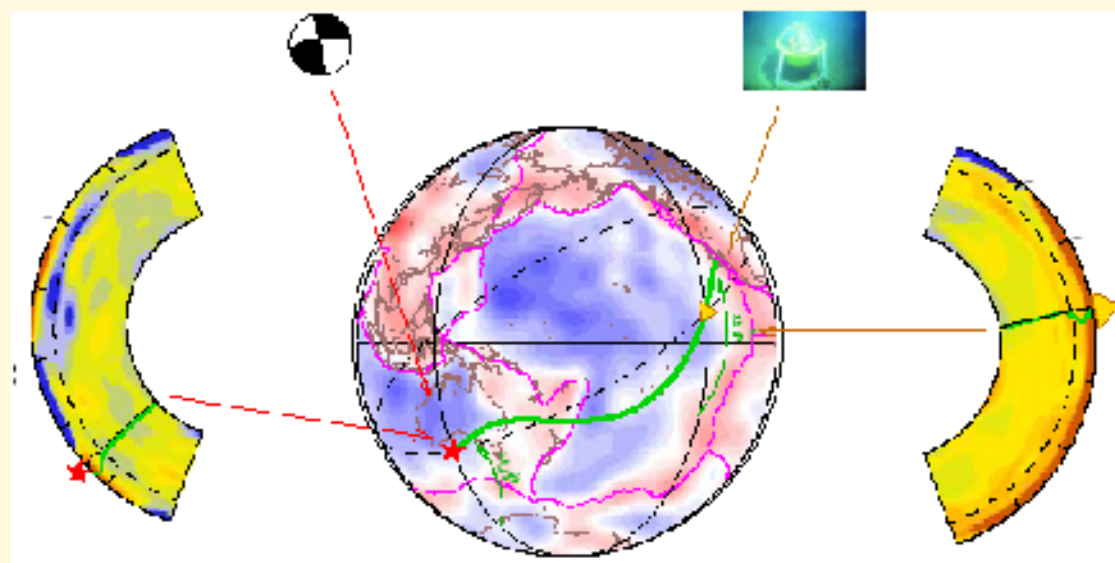


Surface wave full ray theory

Waves locally as plane waves $\lambda \ll \Lambda_\Omega$

$$u(\omega) = \sum_{\text{branches}} \sum_{\text{orbits}} A(\omega) \exp^{-i\phi(\omega)}$$

$$A = A_s A_p A_r \quad \phi = \phi_s + \phi_p + \phi_r$$



Setting up the linearized inverse problem

Phase $\phi - \phi_0 = \omega t - \Delta/v$

Amplitude: source, focalisation, scattering, atténuation

from ray theory(*):
$$T = \int_{\text{ray path}} \frac{1}{v(\mathbf{r}(s))} ds$$

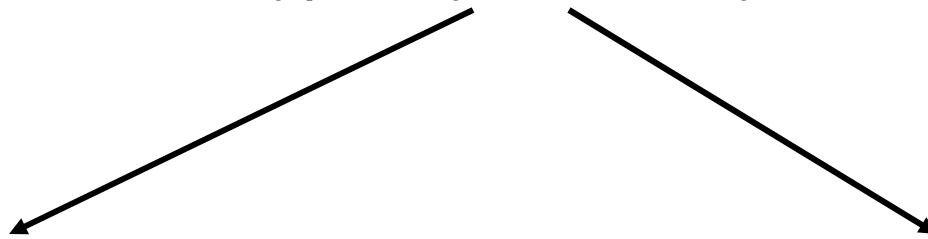
Least-squares solution

given a linear inverse problem: $\delta\mathbf{T} = \mathbf{A} \cdot \mathbf{c}$.

its least-squares solution:

$$\mathbf{c}_{LS} = (\mathbf{A}^T \cdot \mathbf{A})^{-1} \cdot \mathbf{A}^T \cdot \delta\mathbf{T}$$

is typically solved by either:



direct algorithms (SVD, Cholesky...)

iterative algorithms (CG, LSQR...)

Setting up the linearized inverse problem

from ray theory:

$$T = \int_{\text{ray path}} \frac{1}{v(\mathbf{r}(s))} ds$$

first-order Taylor expansion: $\delta T = - \int_{\text{ray path}} \frac{\delta v(\mathbf{r})}{v_0^2(\mathbf{r}(s))} ds$

“parameterization”:

$$\delta v(\mathbf{r}) = \sum_{i=1}^N c_i f_i(\mathbf{r})$$

Unknown
coefficients are
constant and don't
need to be

$$\delta T = - \sum_{i=1}^N c_i \int_{\text{ray path}} \frac{f_i(\mathbf{r})}{v_0^2(\mathbf{r}(s))} ds$$

Setting up the linearized inverse problem

$$\delta T = - \sum_{i=1}^N c_i \int_{\text{ray path}} \frac{f_i(\mathbf{r})}{v_0^2(\mathbf{r}(s))} ds$$

The same eq. Applies for many data and corresp. ray paths:

$$\delta T_j = - \sum_{i=1}^N c_i \int_{\text{ray path}_j} \frac{f_i(\mathbf{r})}{v_0^2(\mathbf{r}(s))} ds \quad (j=1, \dots, M)$$

given a linear inverse problem:

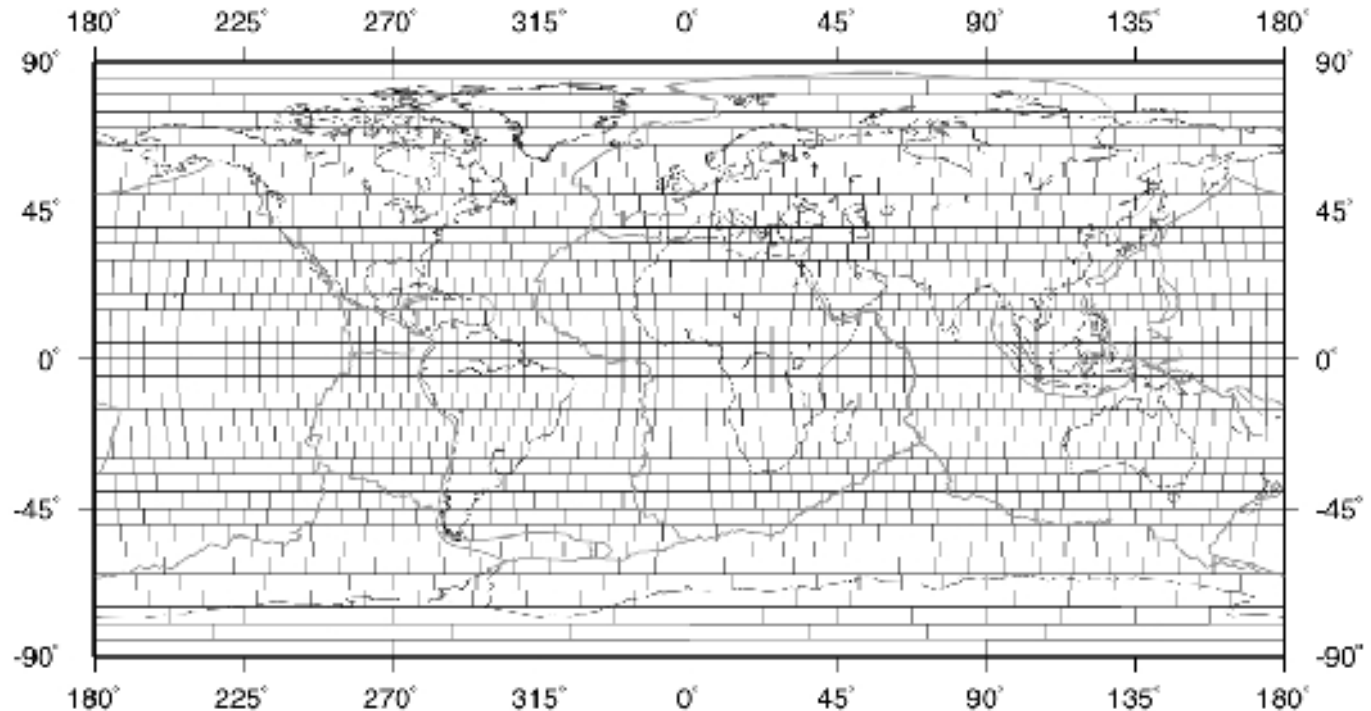
$$\delta \mathbf{T} = \mathbf{G} \delta \mathbf{v}$$

$$\delta \mathbf{T} = \mathbf{A} \cdot \mathbf{c}$$

Parameterization, or choice of basis functions

$$\delta v(\mathbf{r}) = \sum_{i=1}^N c_i f_i(\mathbf{r})$$

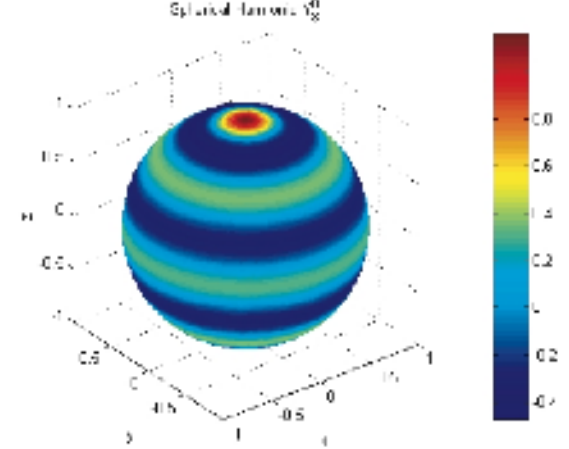
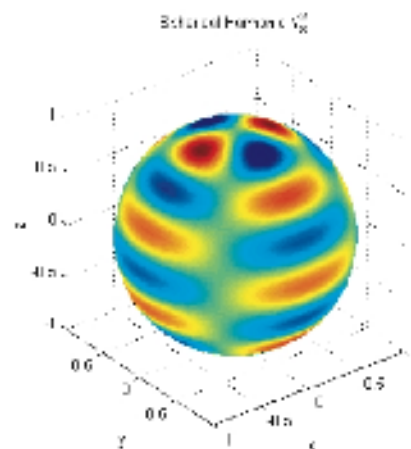
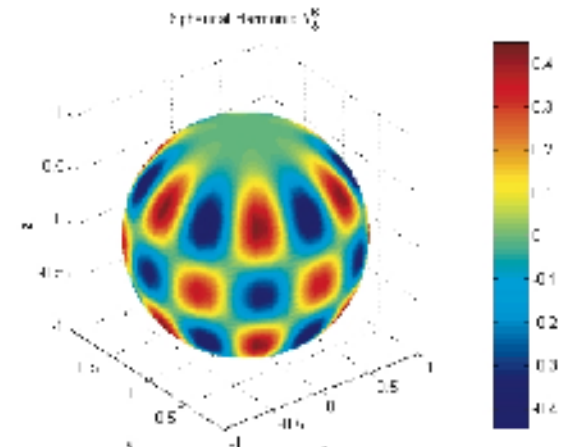
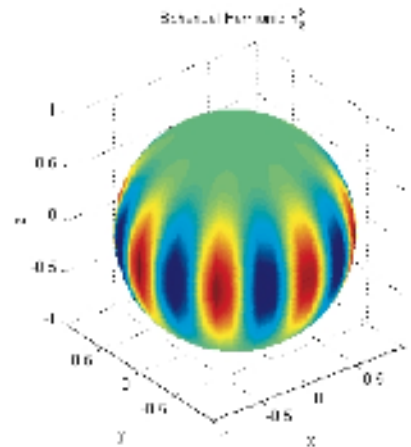
pixels (“local”
basis
functions)



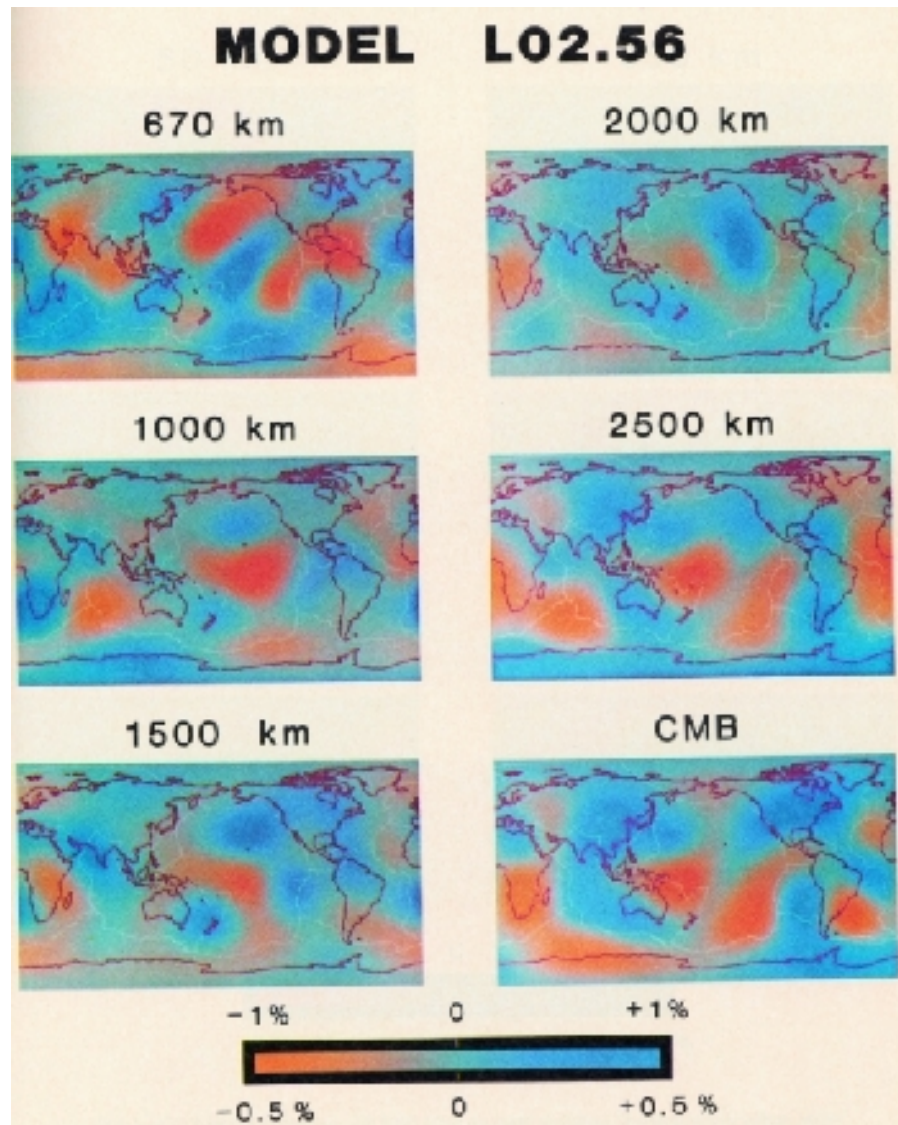
Parameterization, or choice of basis functions

$$\delta v(\mathbf{r}) = \sum_{i=1}^N c_i f_i(\mathbf{r})$$

spherical harmonics
("global" basis
functions)

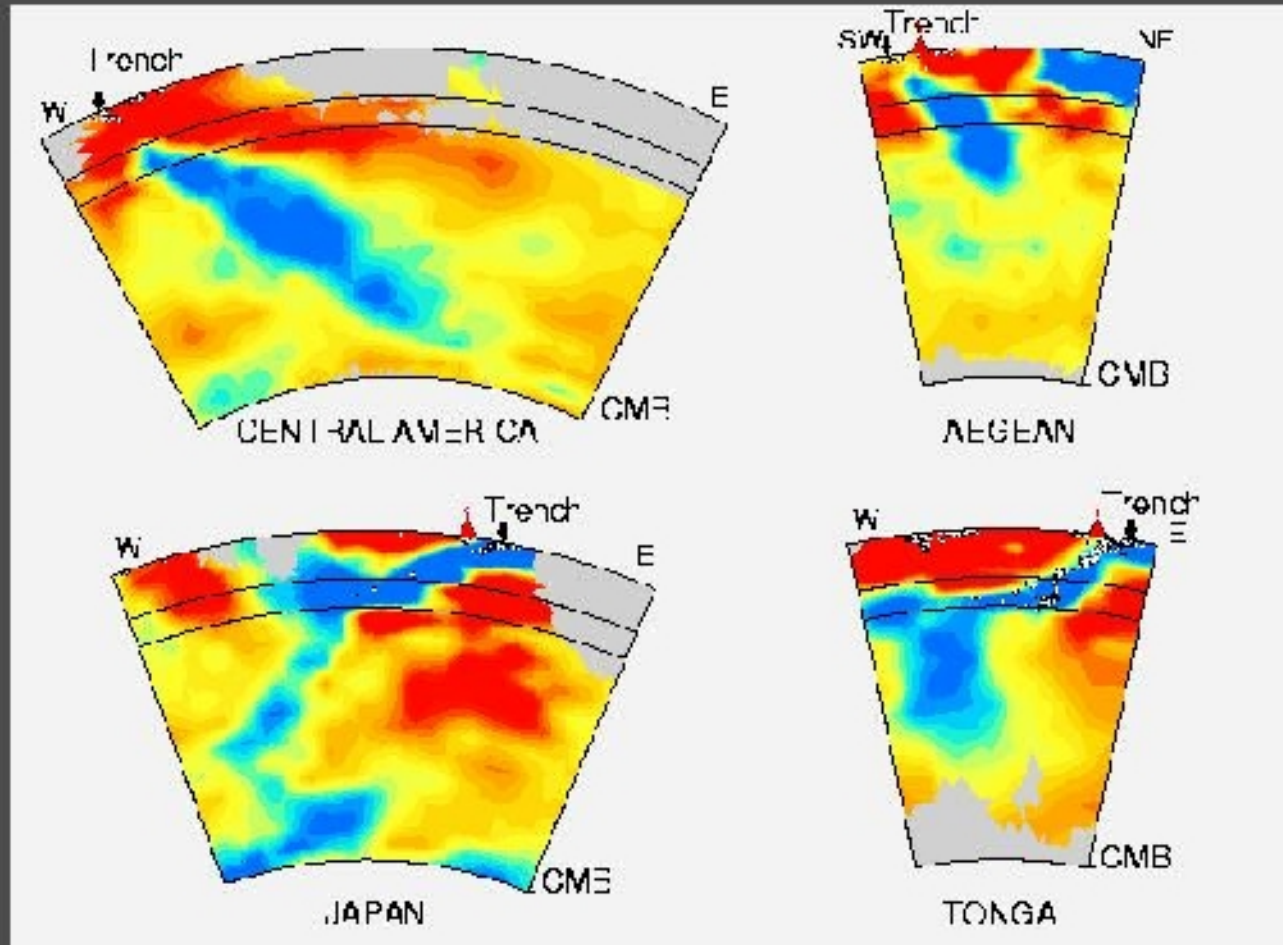


body wave tomography: examples

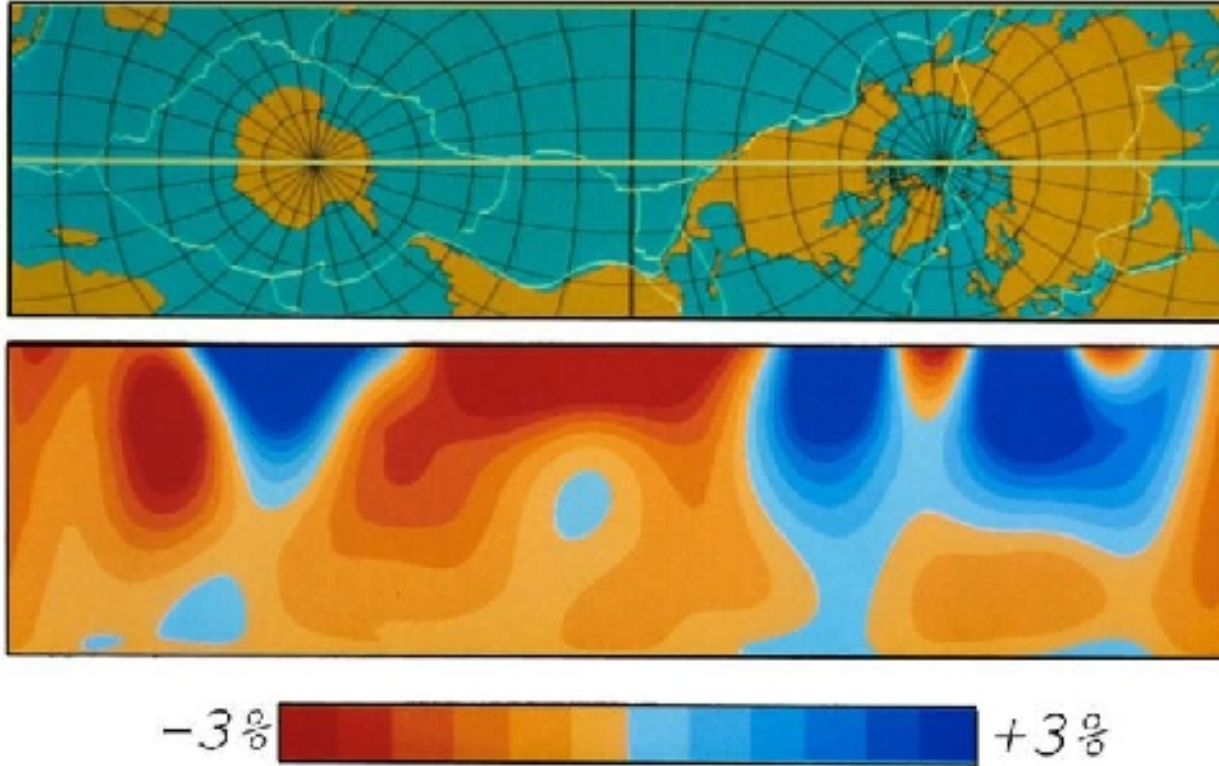


1980s low-degree model (Dziewonski)

Imaging of Slabs



body wave (+surface wave) tomography: examples



1980s low-degree model (Dziewonski & Woodhouse), cross-section

Tomographie d'onde de surface

Phase $\phi - \phi_0 = \omega t - \Delta/c$

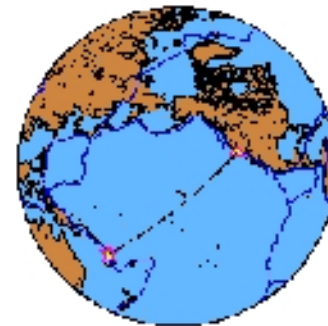
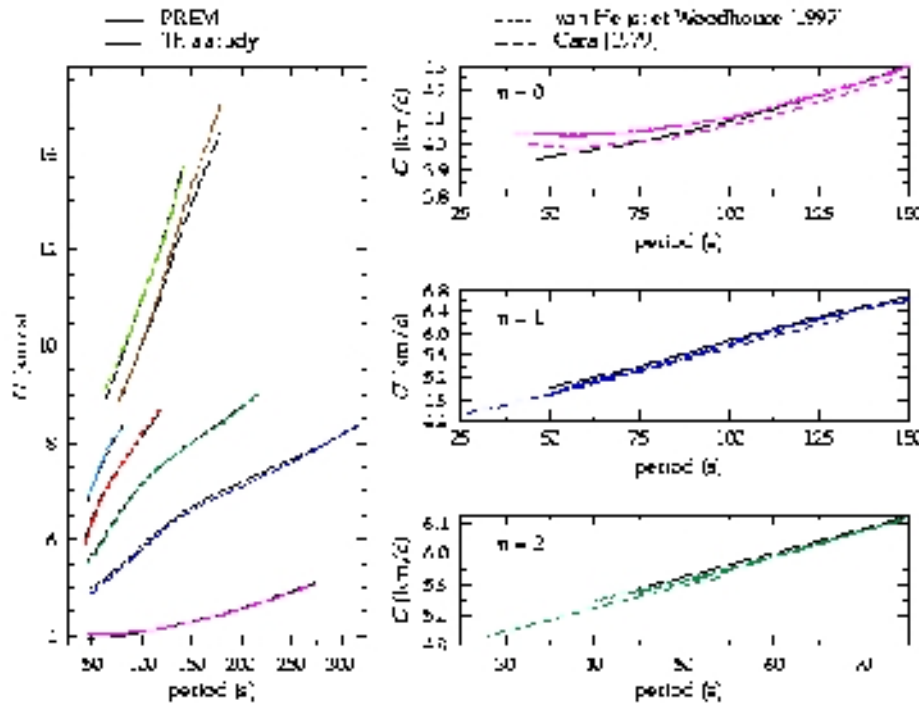
$C(T)$ phase velocity, $C_0(T)$ ref. phase velocity

What tomographers need is the relationship between displacement \mathbf{u} and Earth's "structure", $\rho(\mathbf{r})$ $\lambda(\mathbf{r})$ $\mu(\mathbf{r})$

we can *linearize* it in a perturbative approach:

$$\delta \mathbf{u} = \int_V [\mathbf{K}_\rho(\mathbf{r}) \delta \rho(\mathbf{r}) + \mathbf{K}_\lambda(\mathbf{r}) \delta \lambda(\mathbf{r}) + \mathbf{K}_\mu(\mathbf{r}) \delta \mu(\mathbf{r})] dV$$

1st step: Calculation of dispersion curves: Fundamental modes and higher modes

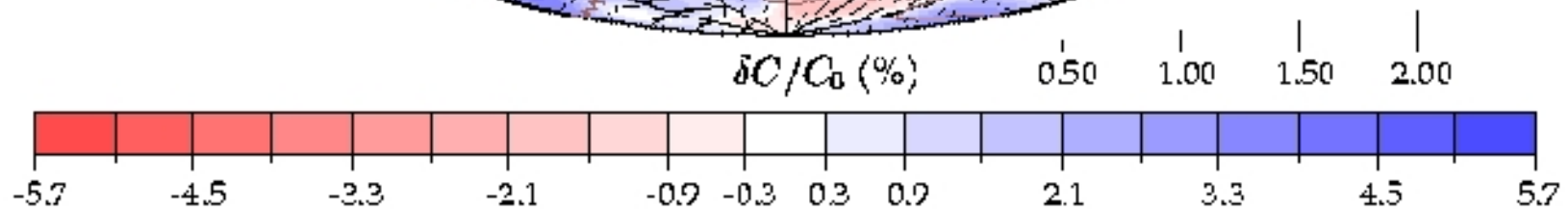
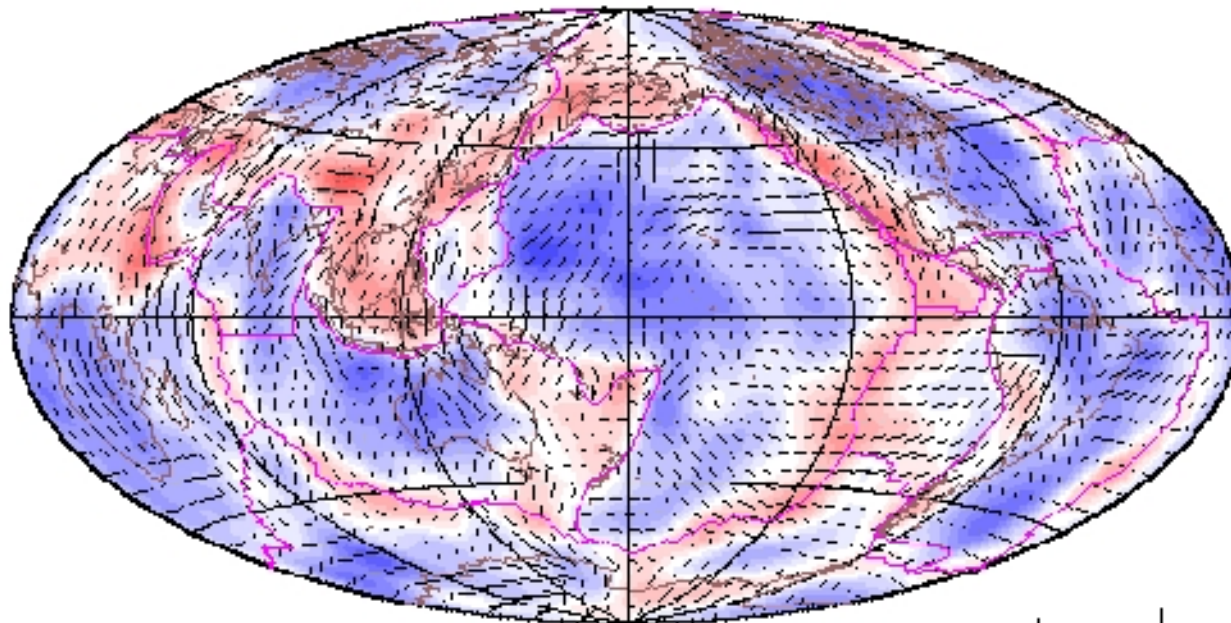


Comparison with previous results along the Vanuatu-California path.

2nd step: regionalization

ÉRIC BEICLER

2002



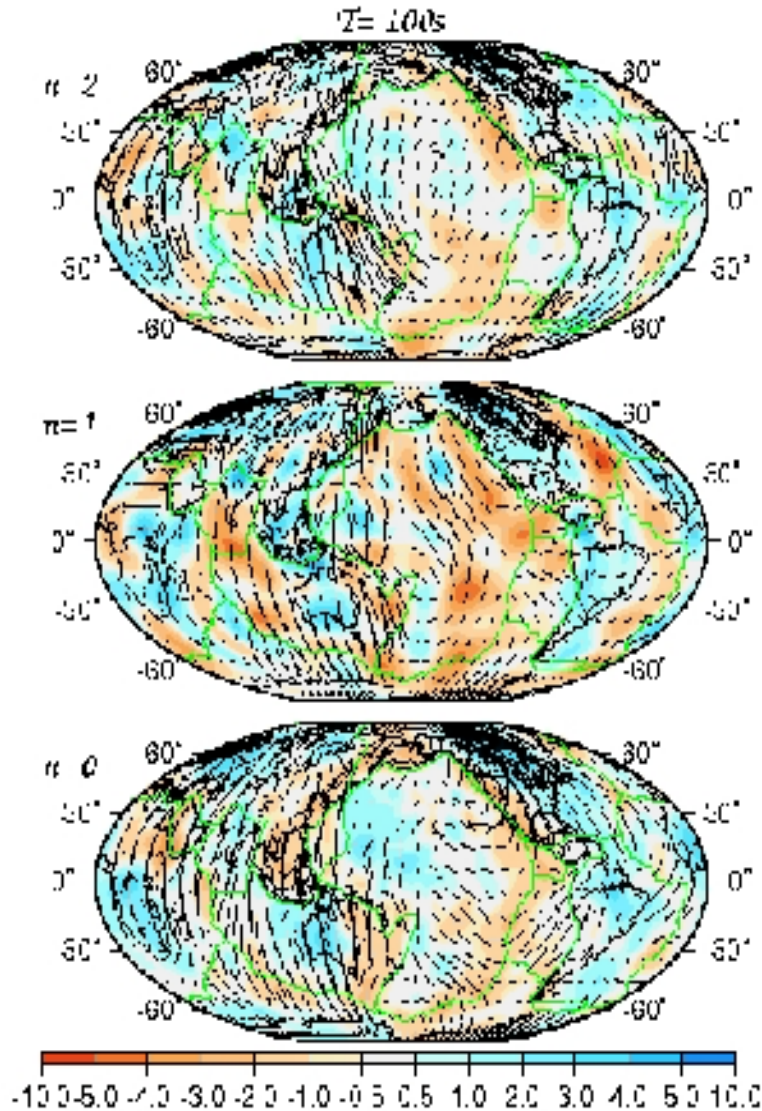
Anisotropic model resulting from phase velocities inversion (C.L.A.S.H., including 2Ψ and 4Ψ terms),
 $n = 0$, $T = 50$ s.

Phase velocity maps At 100s

2nd overtone

1st overtone

Fundamental mode



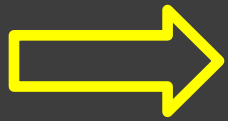
Global Tomography

3rd step: Inversion at depth

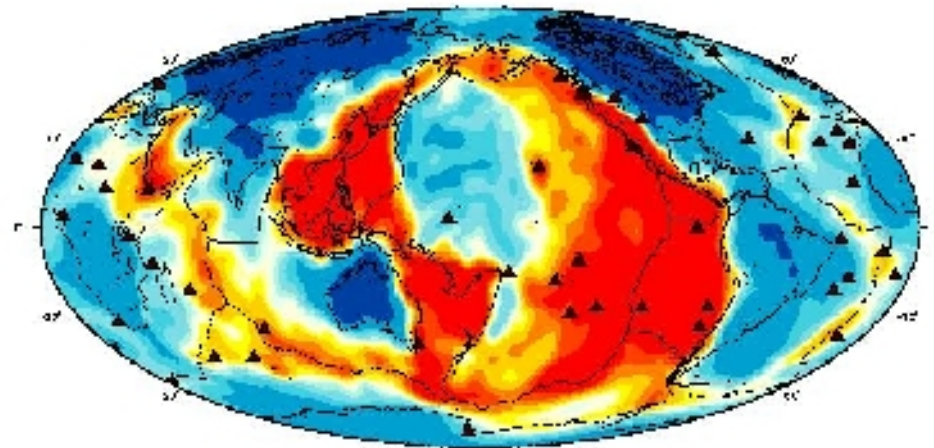
Scale $\Lambda \approx 2000\text{km}$ (degree 20)

Seismic wavelength $\lambda \leq 500\text{km}$

Ray theory applies



Shear wave velocities - depth = 100km



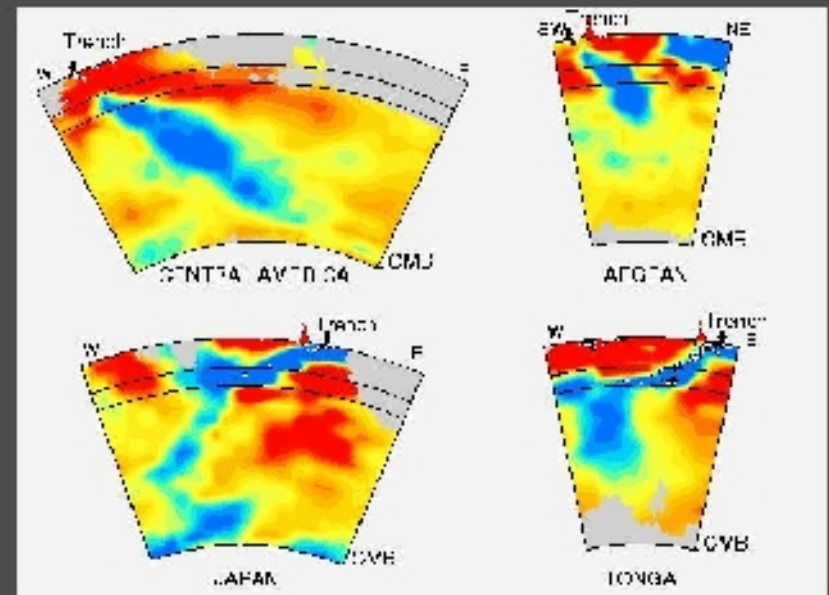
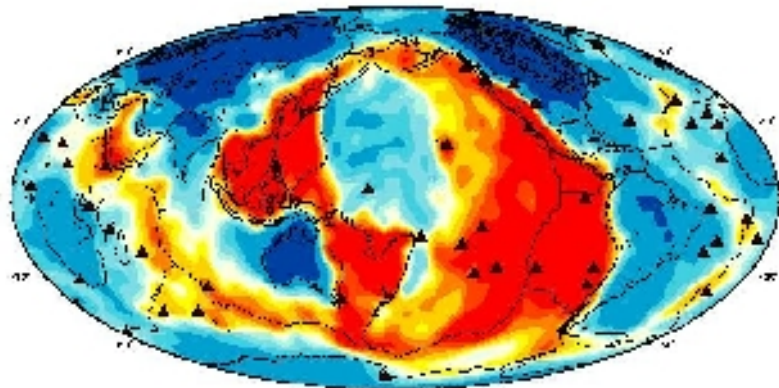
Global and Regional Tomography

From Global Scale to Regional scale

Scale $\Lambda \approx 200\text{-}500\text{km}$

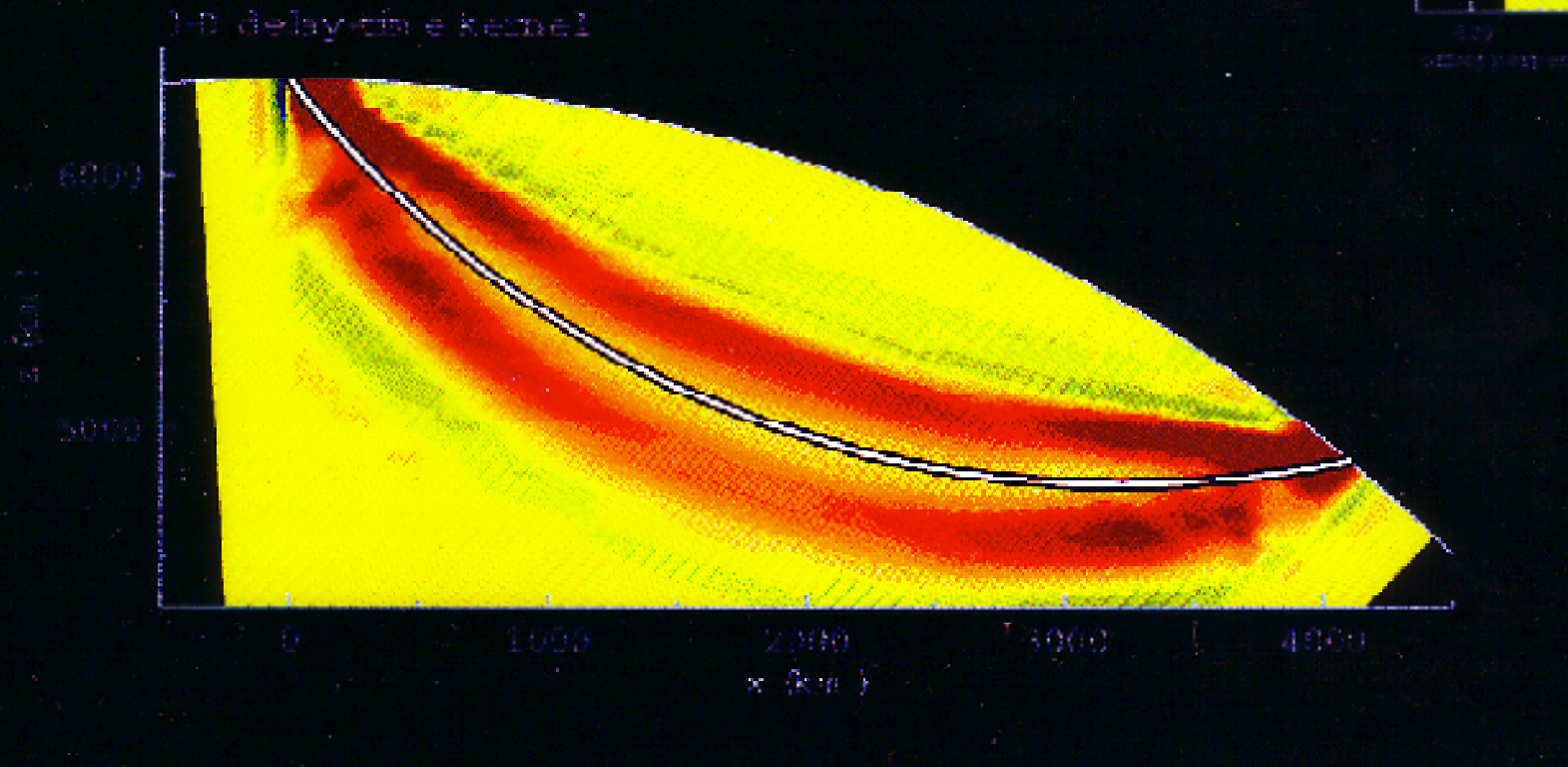
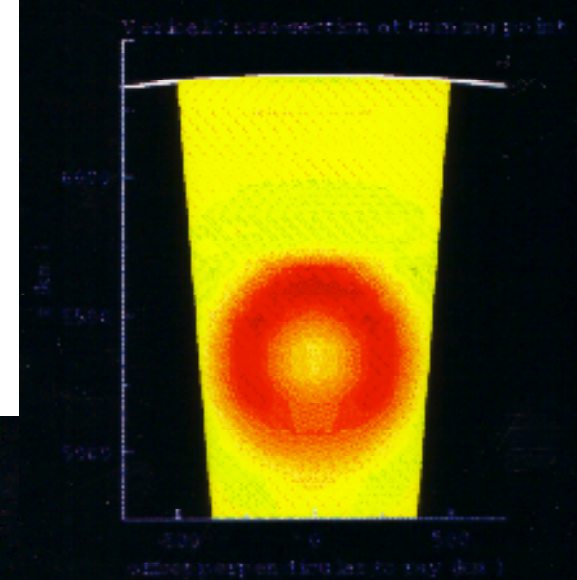
Seismic wavelength $20\text{km} \leq \lambda \leq 500\text{km}$

Anomalies de Vitesses d'ondes S - Profondeur - 1000 kpa



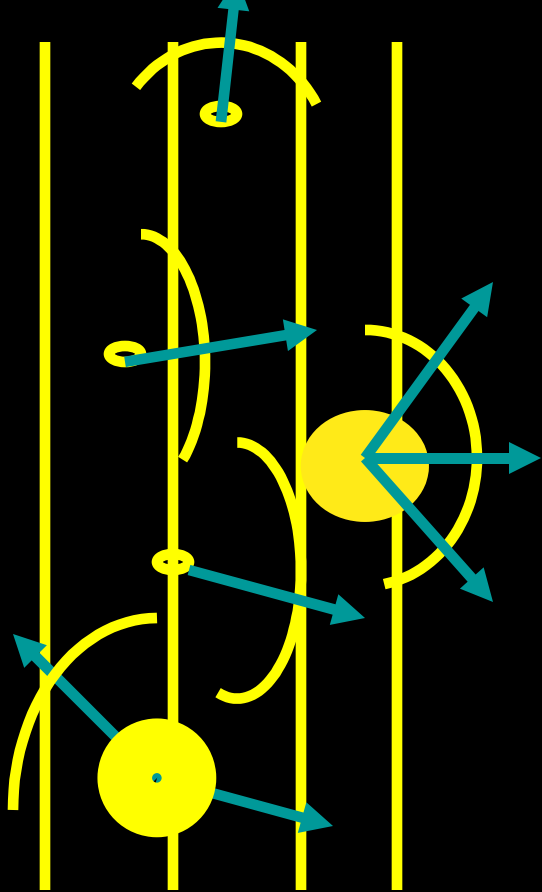
Van der Hilst et al., 1998

Sensitivity of the delay time to the local seismic velocity



Λ heterogeneity scale, λ wavelength

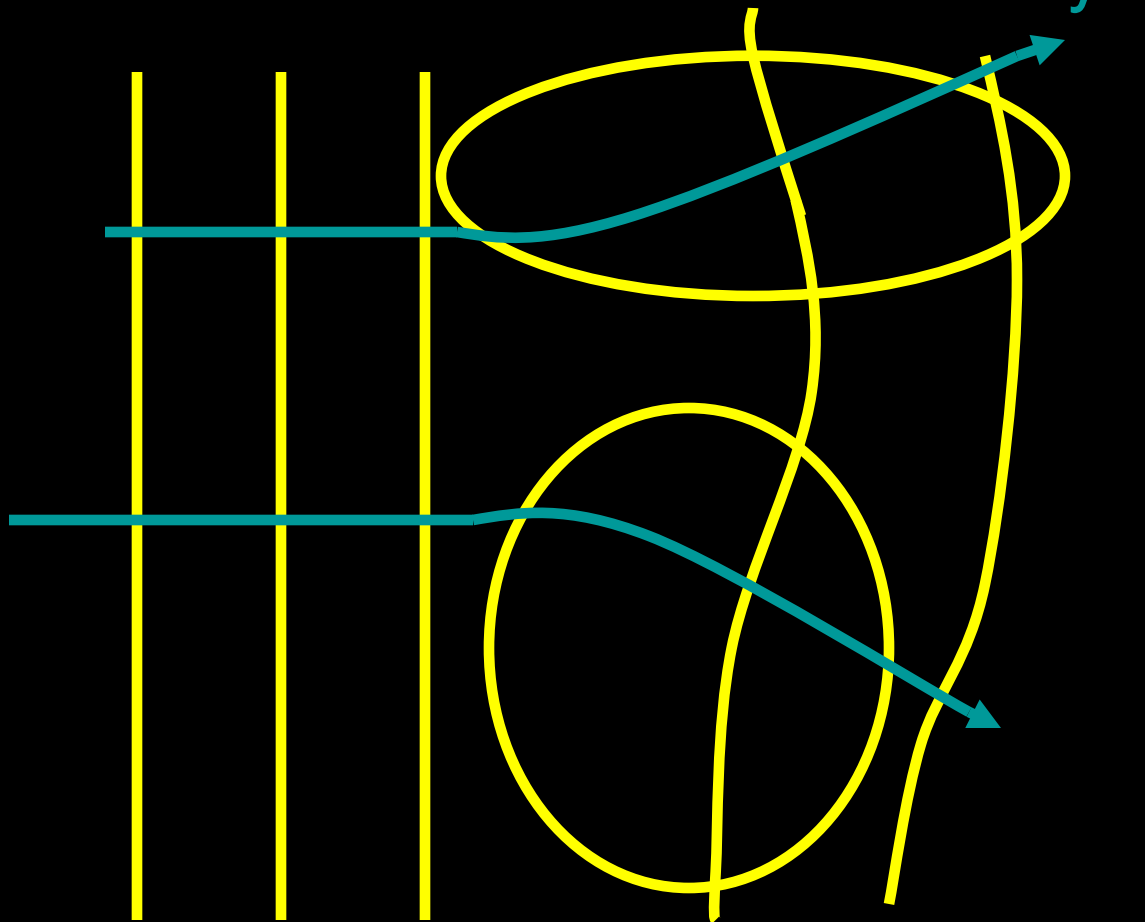
Diffracted waves



wavefronts

$\lambda \sim$

Λ or $\lambda \gg \Lambda$



wavefronts

rays

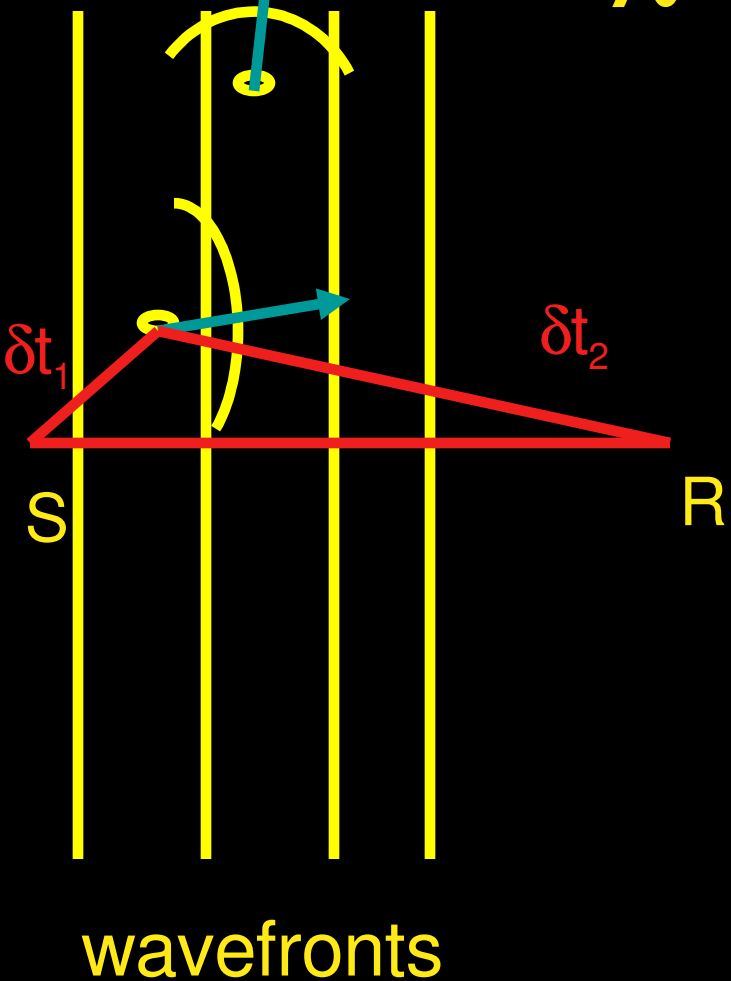
$\lambda \ll \Lambda$

Λ heterogeneity scale, λ wavelength

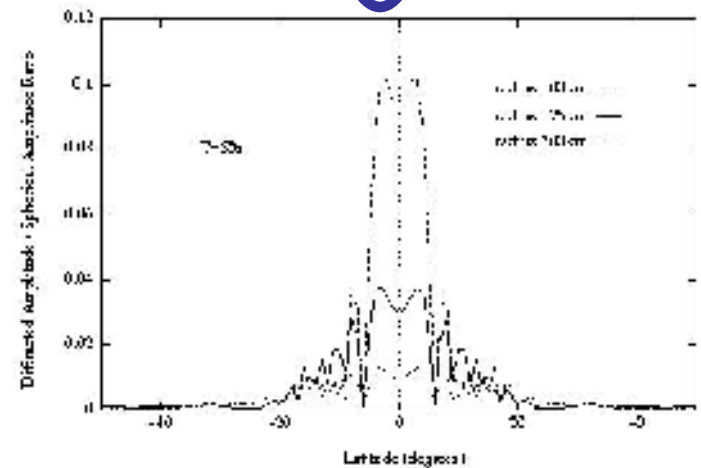
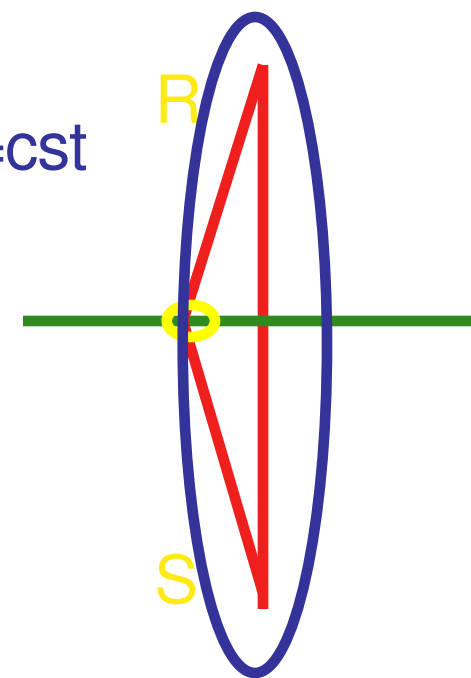
Diffracted waves

$\lambda \sim$

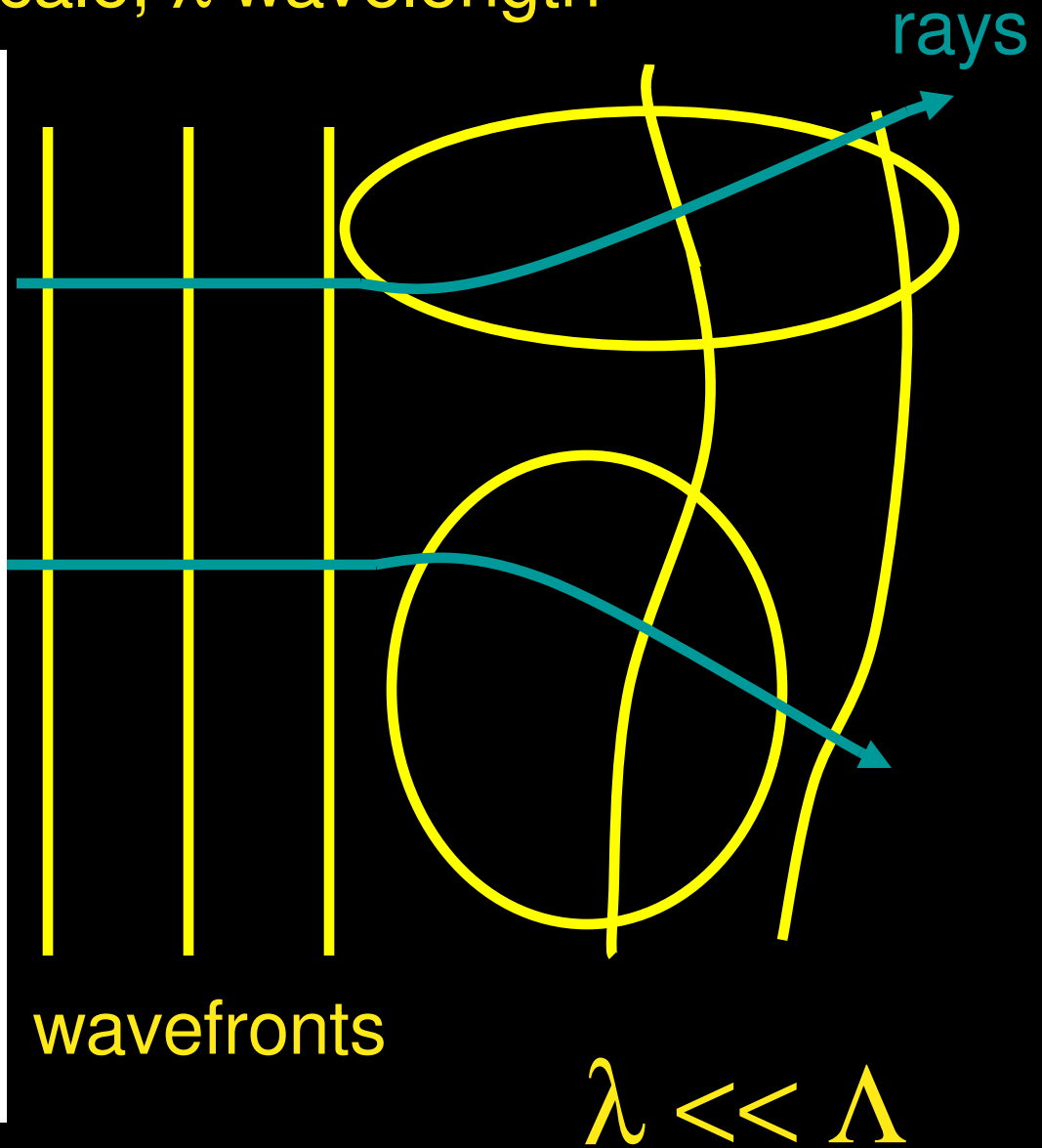
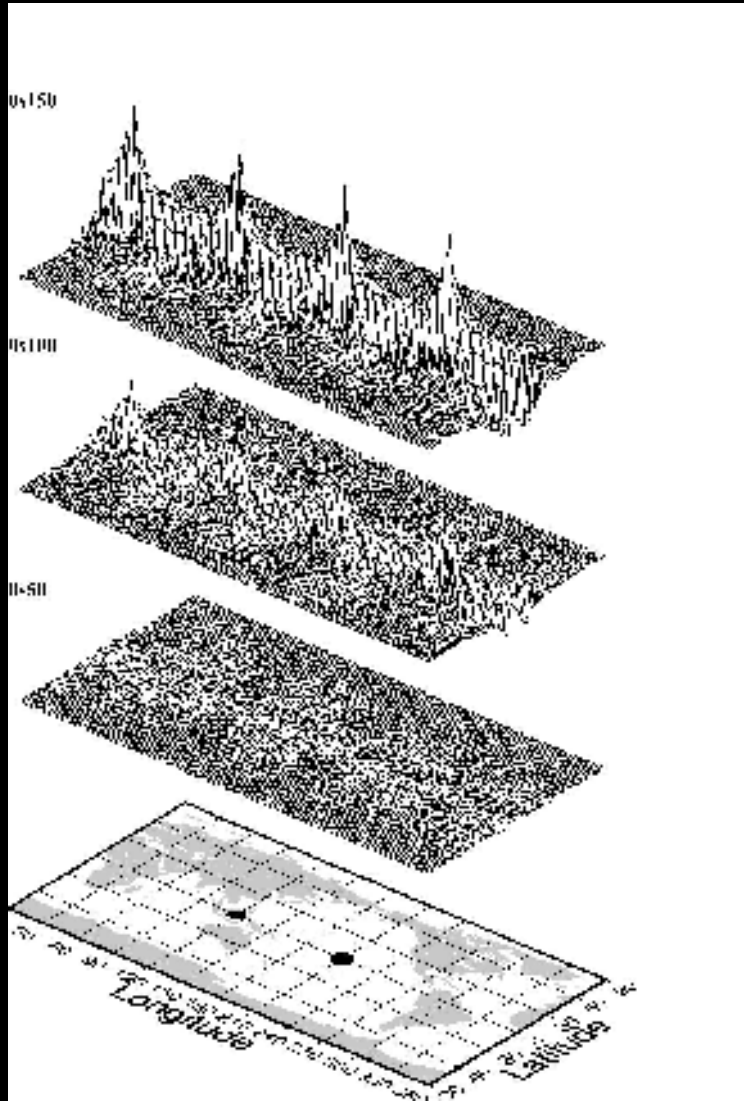
Λ or $\lambda \gg \Lambda$



$$\delta t_1 + \delta t_2 = \text{cst}$$



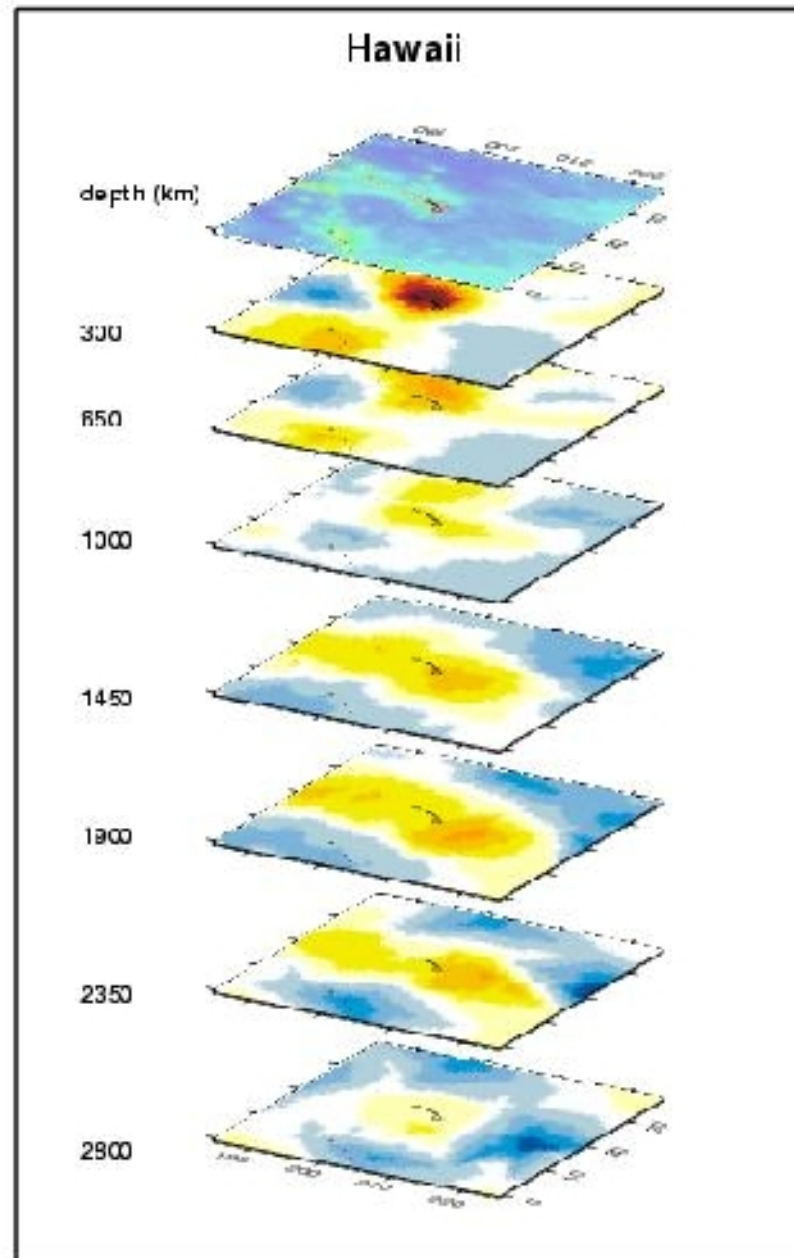
Λ heterogeneity scale, λ wavelength



Typical scales $\Lambda \approx 2,000\text{km}$, $\lambda \approx 500\text{km}$

Banana-Doughnut Theory (Dahlen et al.)

Application to
global
tomography
(Montelli et al.,
Science, 2004)

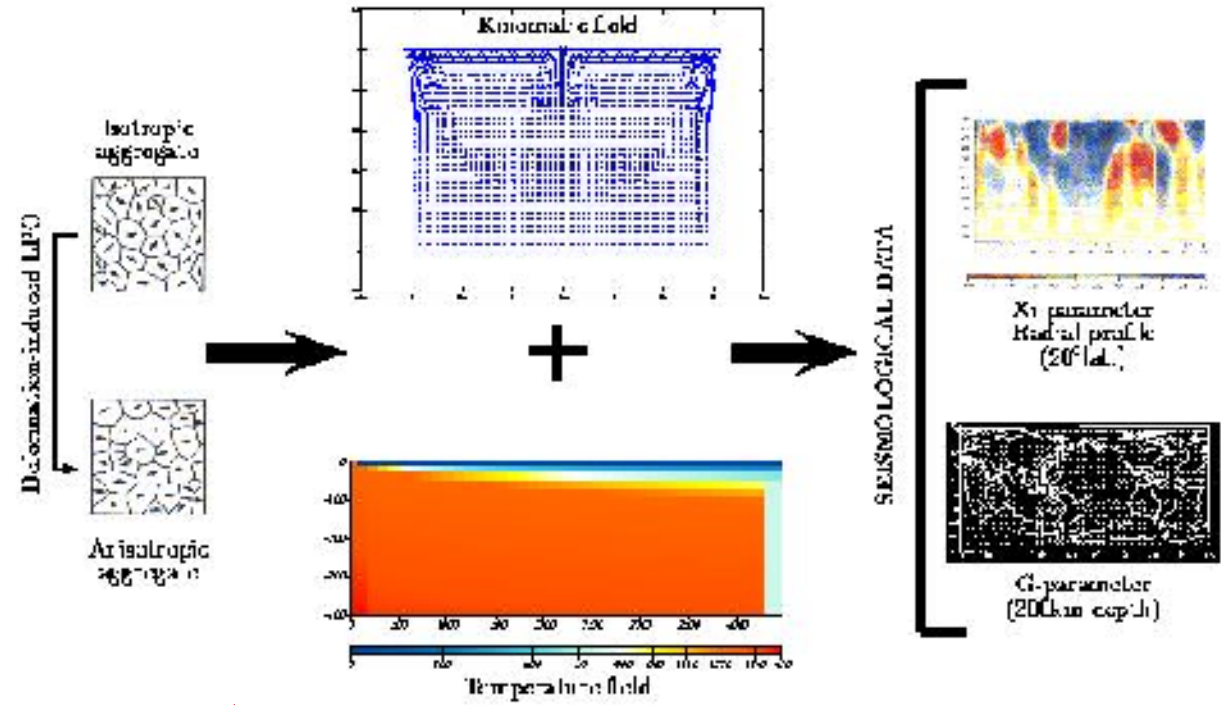




Importance of seismic anisotropy

ANISOTROPY is the Rule not the Exception

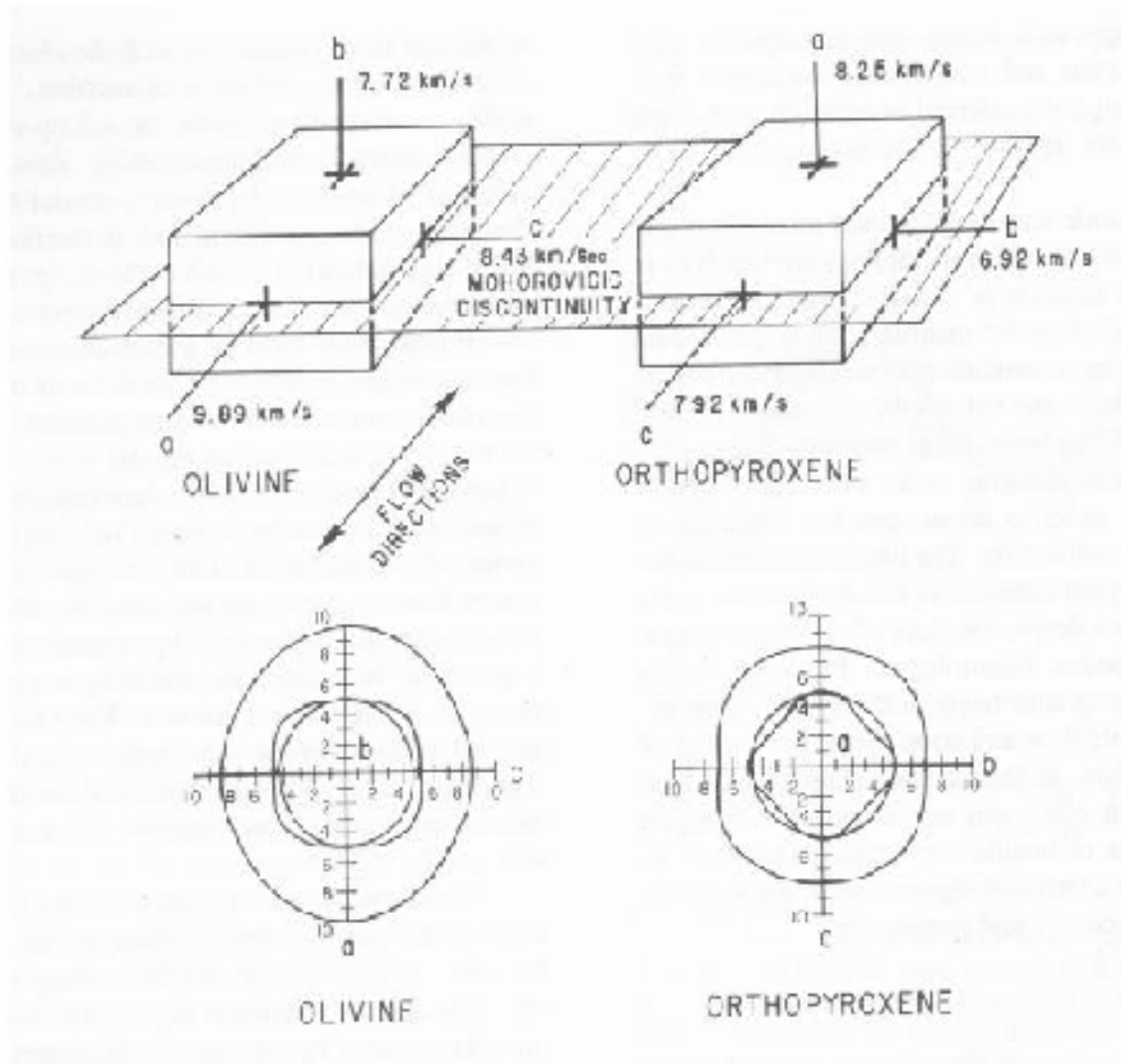
Seismic Anisotropy is present at all scales



Mineralogical composition

Mapping Convection

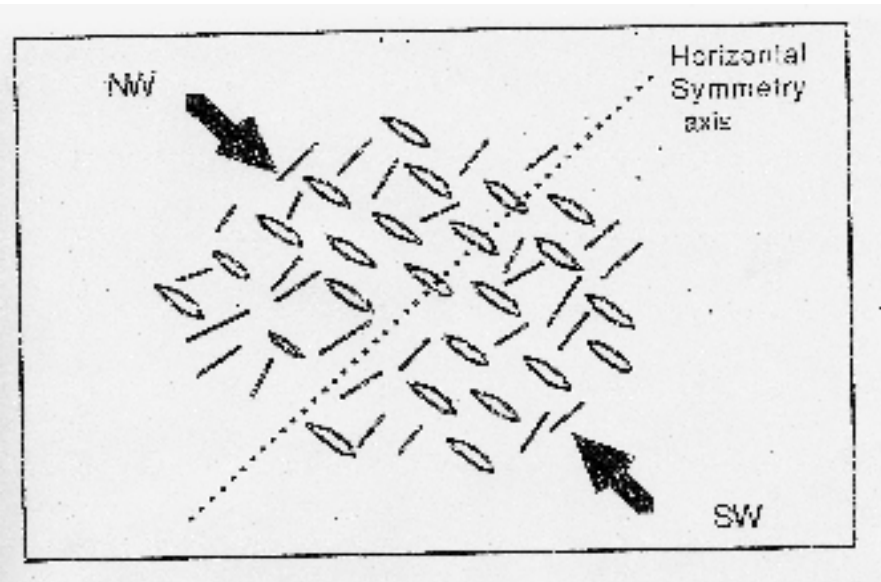
(Montagner and Guillot, 2001)



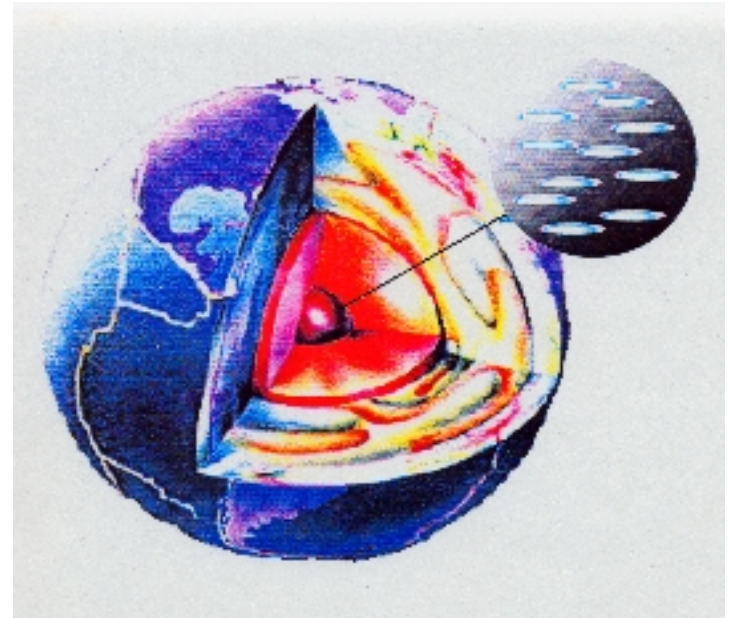
From Christensen and Lundquist, 1982

Cracks, fluid inclusions

Crust



Inner core



(Babuska and Cara, 1991)

(Singh et al., 2001)

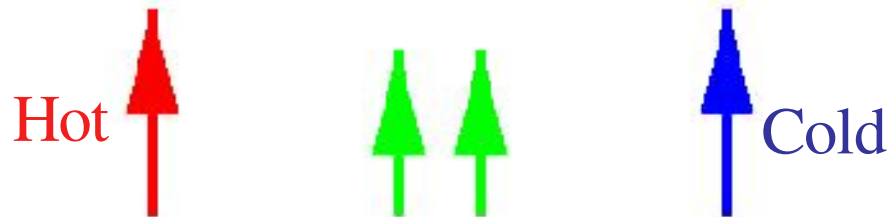
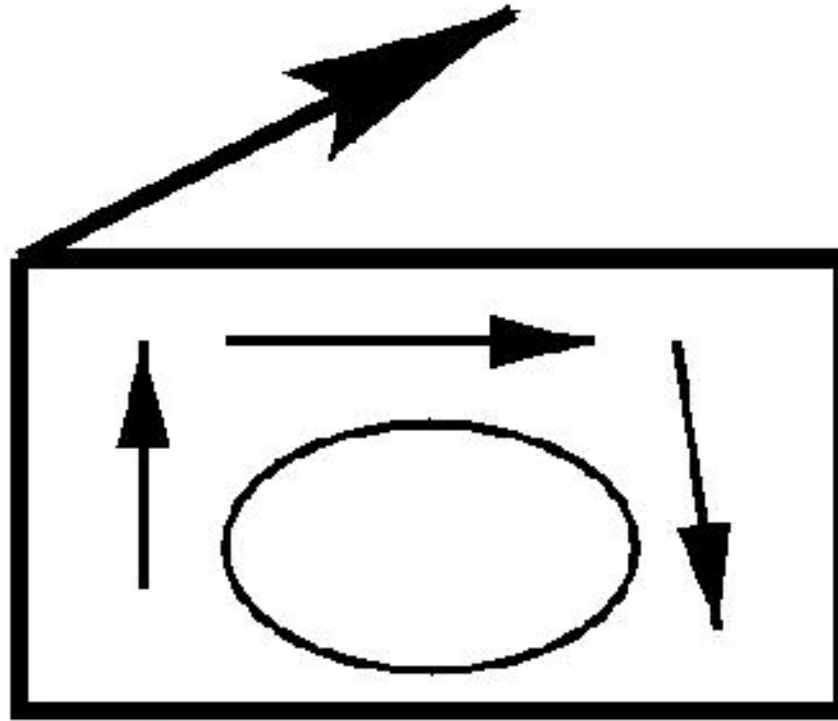


Importance of seismic anisotropy ANISOTROPY is the Rule not the Exception

Anisotropy is present at all scales

- From microscopic scale up to macroscopic scale
- Efficient mechanisms of alignment
(L.P.O.: lattice preferred orientation
S.P.O.: shape preferred orientation; fine layering)

Montagner & Guillot, 2001



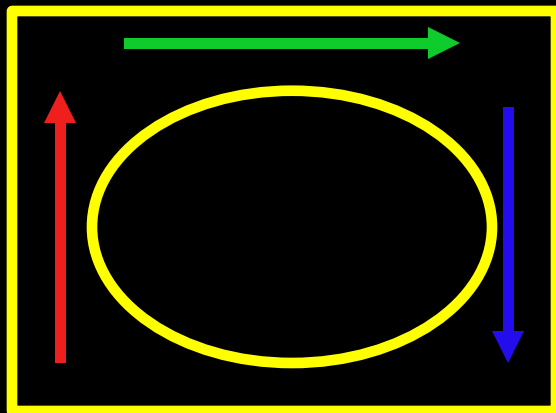
$\Delta\alpha$ Effect of Mineral Orientation

ΔT Effect of Temperature Heterogeneities

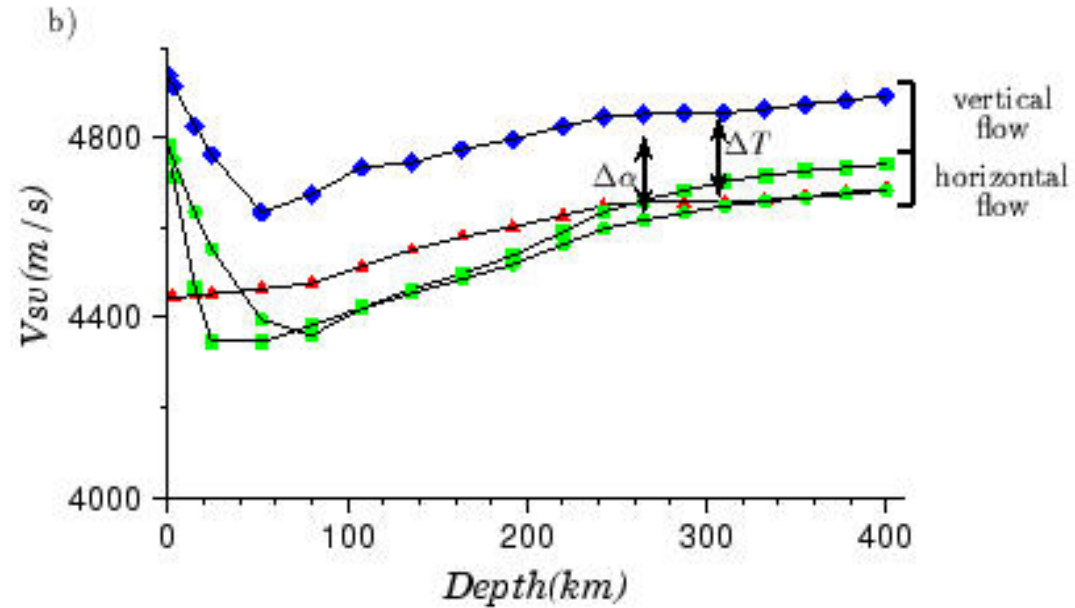
$\Delta\alpha$: Anisotropy
Effect

ΔT : Temperature
Effect

$$\Delta\alpha \approx \Delta T$$



Olivine (60%) + Opx (40%)

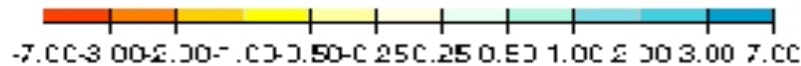
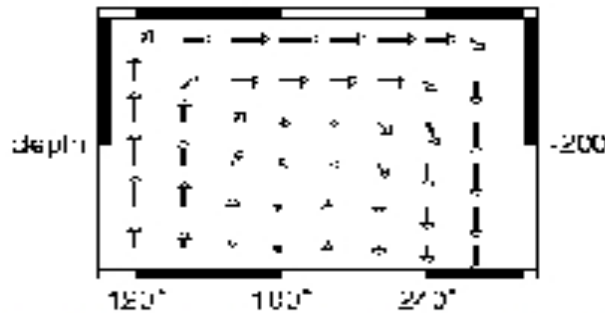
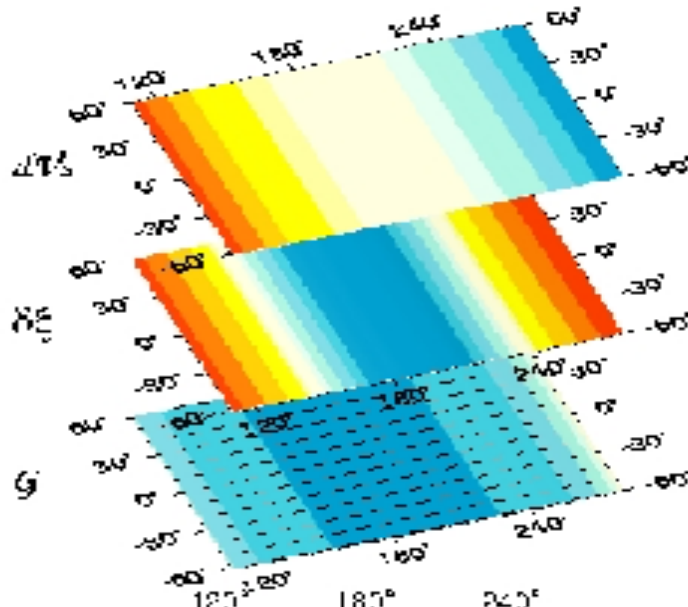


Montagner & Guillot, 2002

S- Velocity

Radial
Anisotropy

Azimuthal
Anisotropy



Interpretation

- Anisotropy is observed on different kinds of seismic waves
- Body waves (Pn, Shear wave splitting)
 - Surface waves (Rayleigh-Love discrepancy; Azimuthal anisotropy)



Importance of seismic anisotropy
ANISOTROPY is the Rule not the Exception



Anisotropy is present at all scales

- from microscopic scale to macroscopic scale
- Efficient mechanisms of alignment
(L.P.O.: lattice preferred orientation
S.P.O.: shape preferred orientation; fine layering)

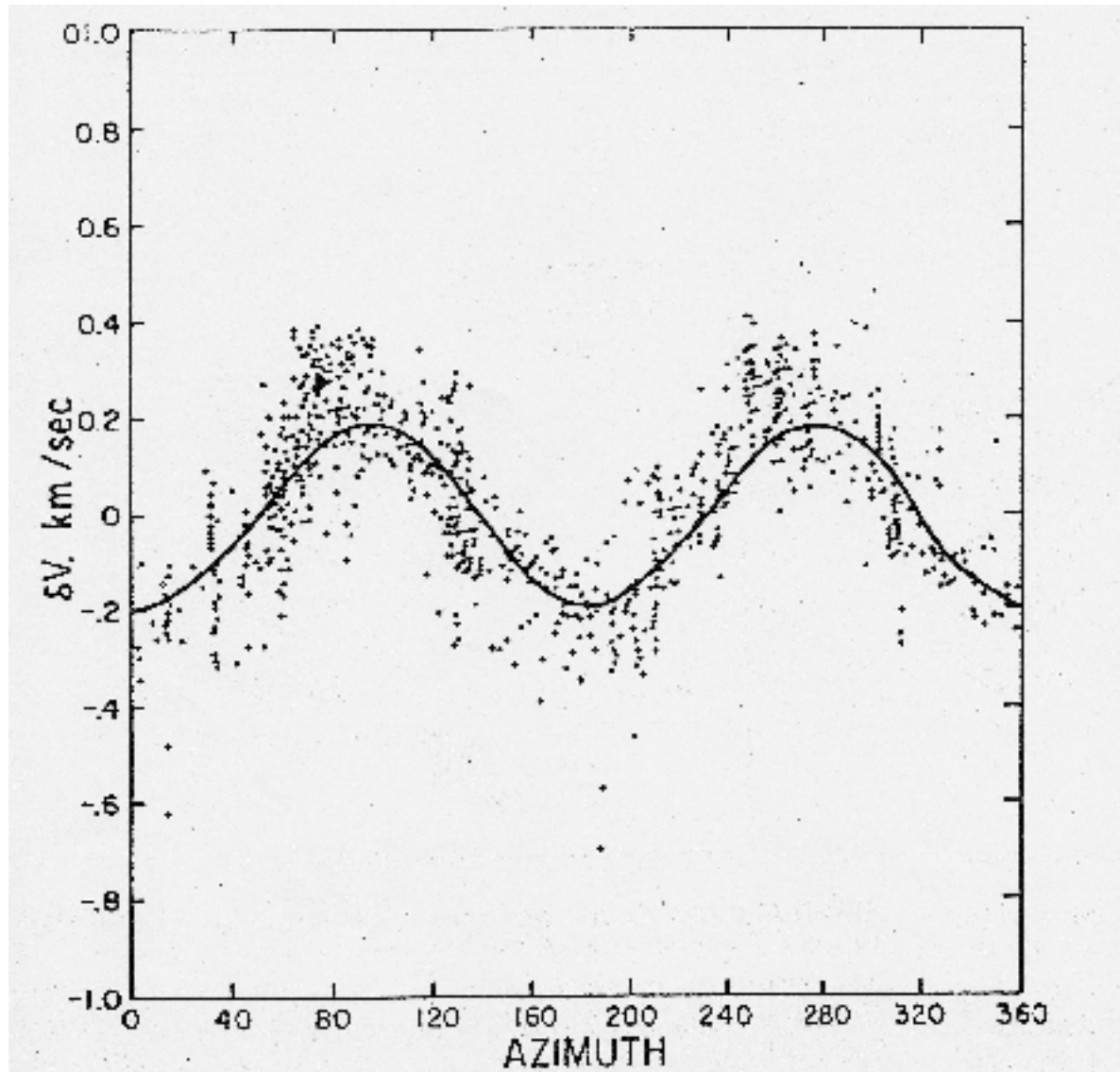
Anisotropy is observed on different kinds of seismic waves

- Body waves (Pn; S-wave splitting)
- Surface waves (Rayleigh-Love discrepancy, azimuthal anisotropy)

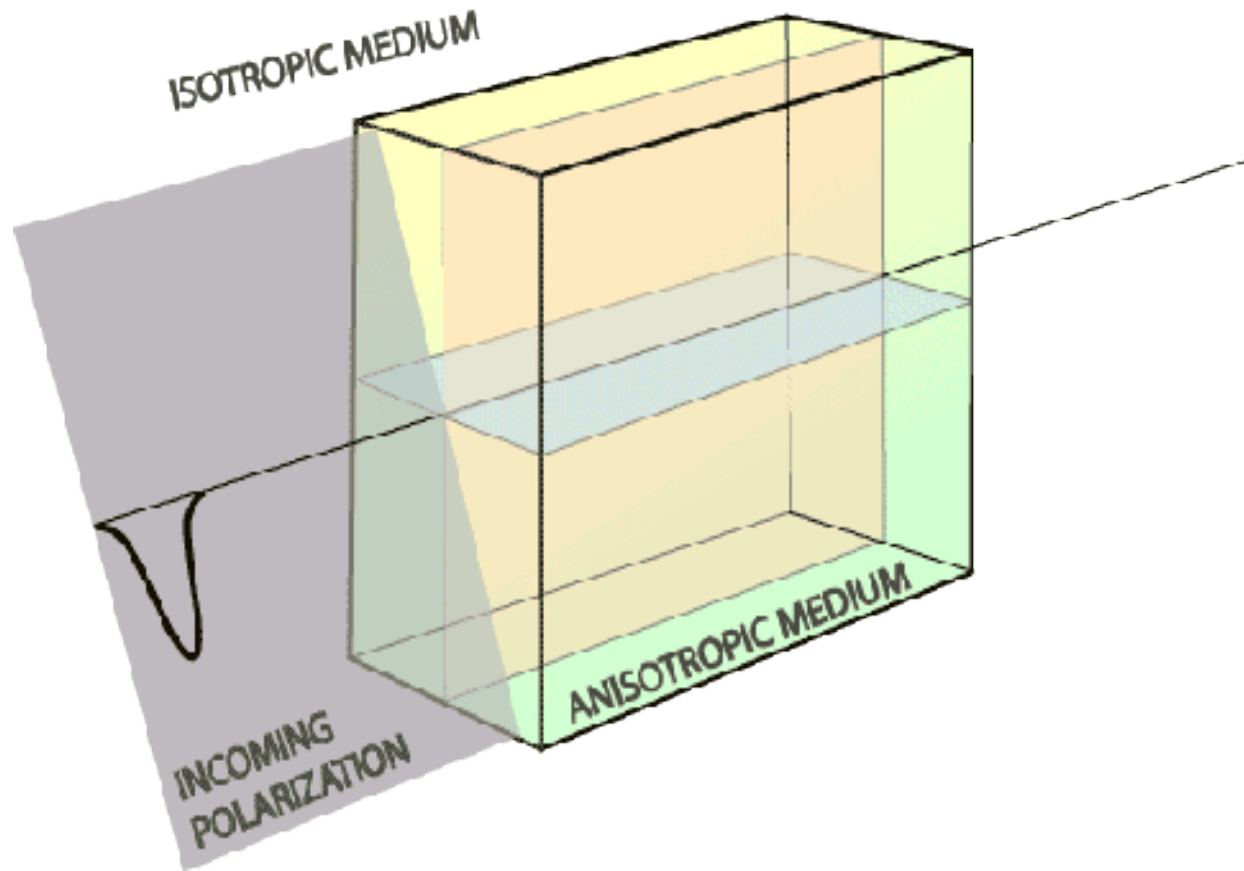
ANISOTROPY REFLECTS AN INNER ORGANIZATION

ANISOTROPY IS NOT A SECOND ORDER EFFECT

Pn- velocities

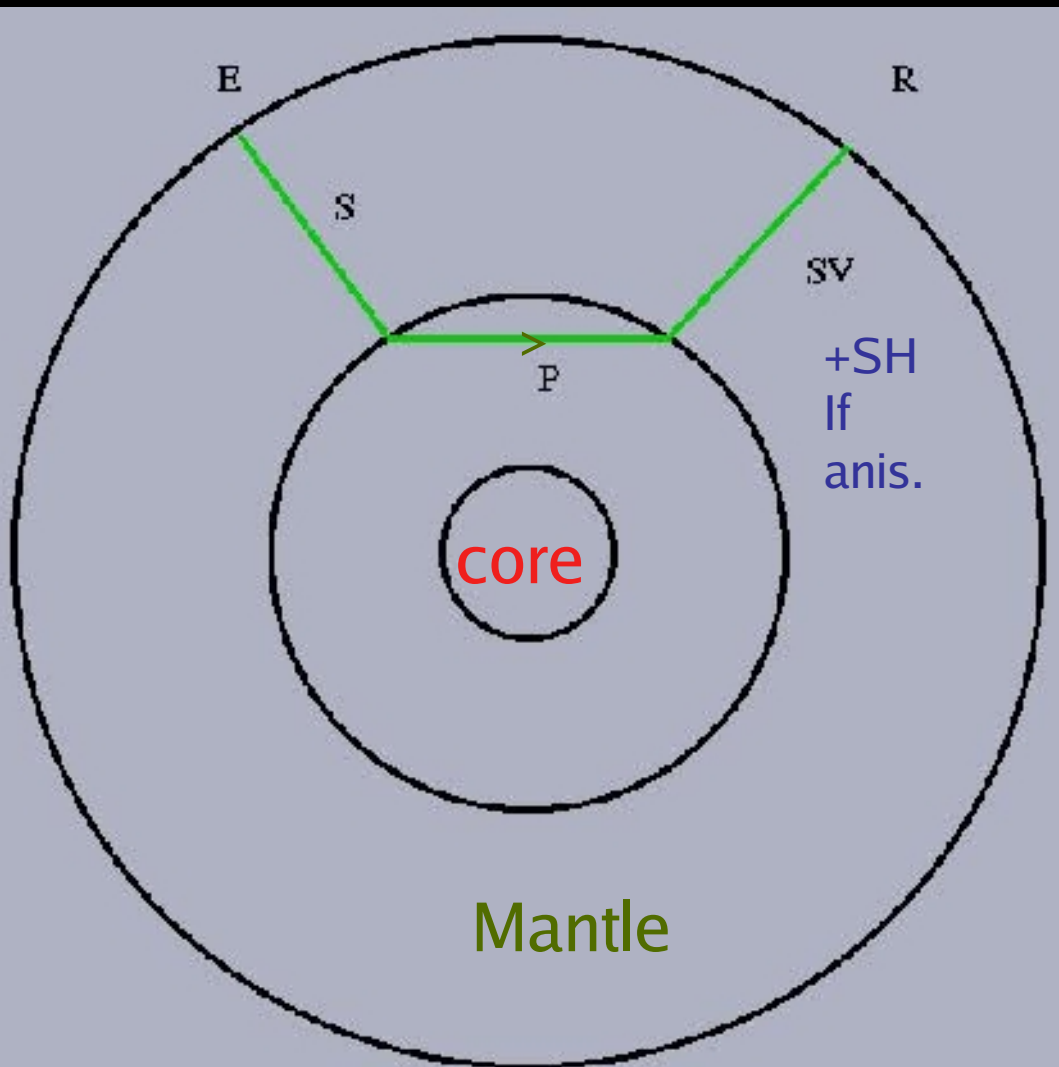


Shear Wave Splitting (Birefringence)

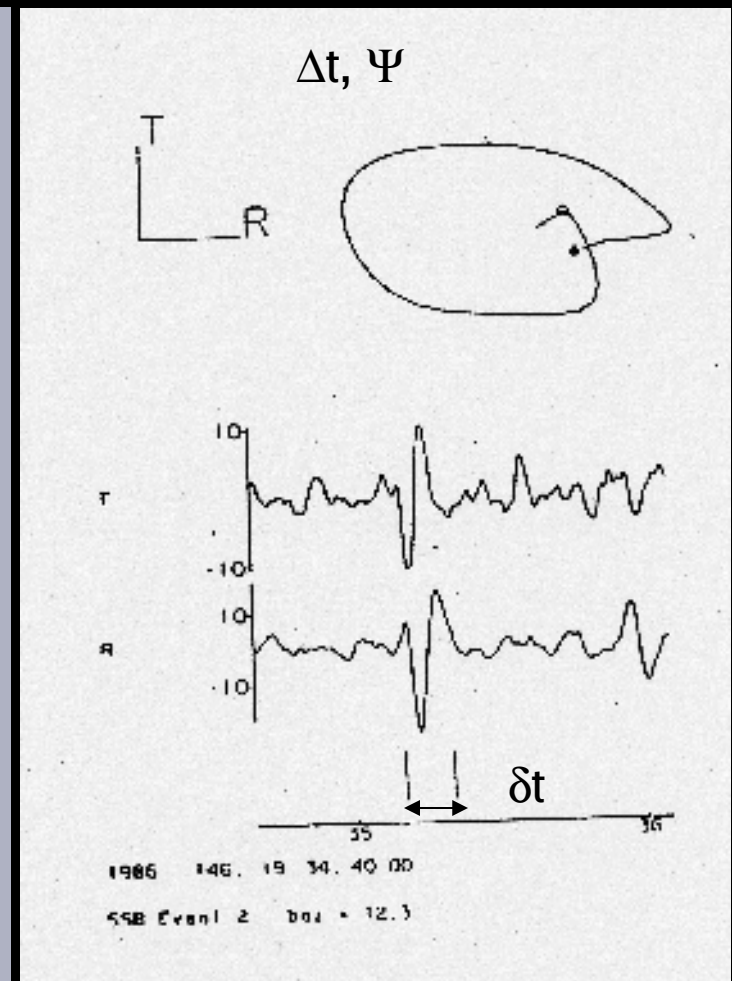


Animation courtesy of Ed Garnero

SKS- Splitting



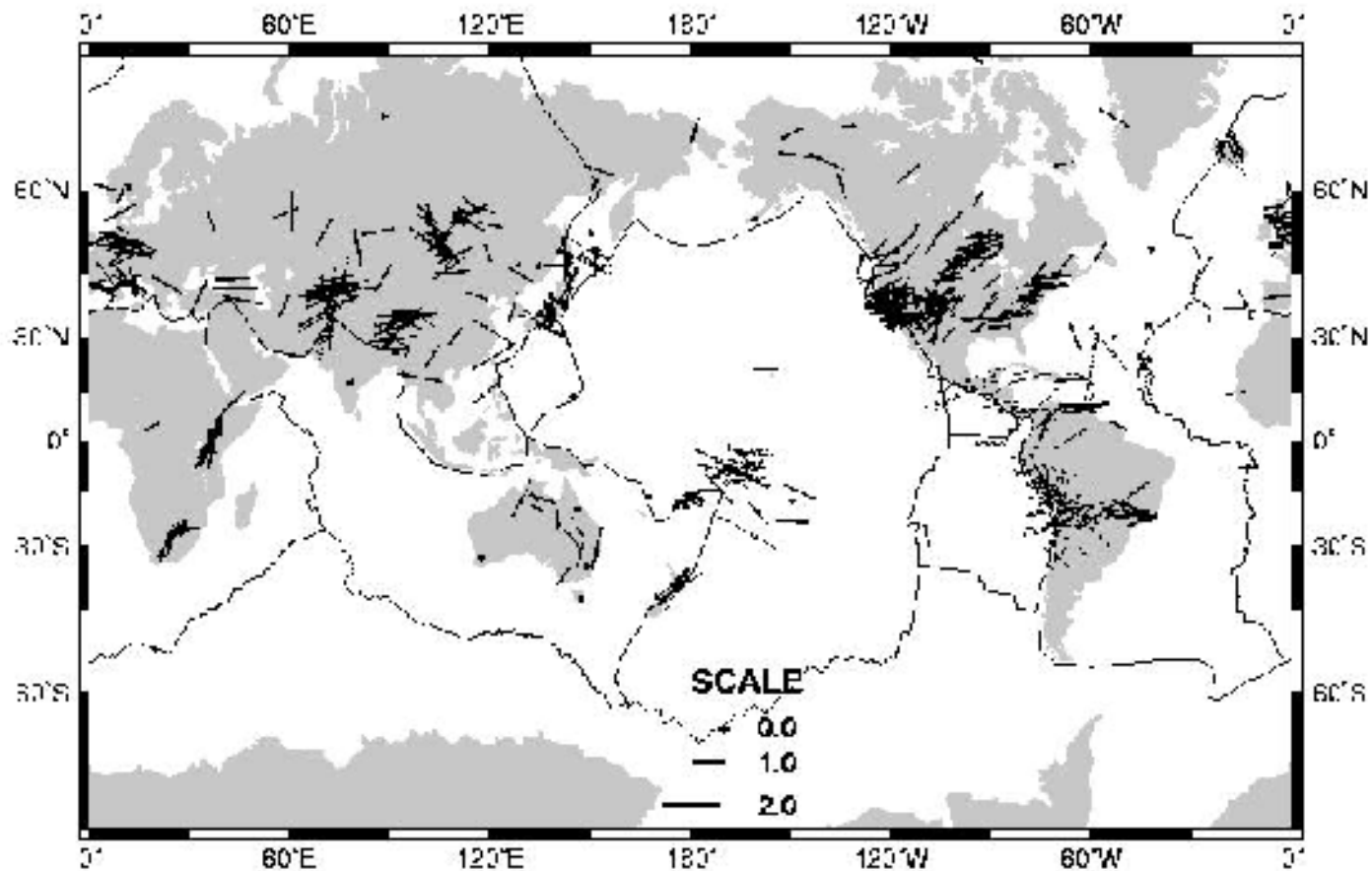
SKS Wave Path



T
R

Vinnik et al., 1989

Compilation of S-wave splitting measurements



Savage, Rev. Geophys., 1999



Importance of seismic anisotropy
ANISOTROPY is the Rule not the Exception



Anisotropy is present at all scales

- from microscopic scale to macroscopic scale
- Efficient mechanisms of alignment
(L.P.O.: lattice preferred orientation
S.P.O.: shape preferred orientation; fine layering)

Anisotropy is observed on different kinds of seismic waves

- Body waves (Pn; S-wave splitting)
- Surface waves (Rayleigh-Love discrepancy, azimuthal anisotropy)

ANISOTROPY REFLECTS AN INNER ORGANIZATION

ANISOTROPY IS NOT A SECOND ORDER EFFECT

Effect of anisotropy on surface waves

Effect on eigenfrequency for multiplet $k=\{n,l,m\}$

$$\frac{\delta\omega_k}{\omega_k} = \frac{\int_{\Omega} \varepsilon_{ij}^* \delta C_{ijkl} \varepsilon_{kl} d\Omega}{\int_{\Omega} \rho_0^* \mathbf{u}_r^* \cdot \mathbf{u}_r d\Omega} \frac{\delta V}{\bar{V}_k}$$

ε strain tensor, u displacement, δC_{ijkl} elastic tensor perturbation

Phase velocity perturbation $V(T,\theta,\phi,\Psi)$ at point $r(\theta,\phi)$ (Smith&Dahlen, 1973)

$$\begin{aligned} \delta V(T,\theta,\phi,\Psi)/V = & \alpha_0(T,\theta,\phi) + \alpha_1(T,\theta,\phi)\cos 2\Psi + \alpha_2(T,\theta,\phi)\sin 2\Psi \\ & + \alpha_3(T,\theta,\phi)\cos 4\Psi + \alpha_4(T,\theta,\phi)\sin 4\Psi \end{aligned}$$

Ψ Azimuth (angle between North and wave vector)

13 parameters

- Best Resolved parameters for Surface Waves

$L = \rho V_{SV}^2$ Isotropic part of V_{SV} .

$\xi = \frac{N}{L} = \frac{V_{SH}^2}{V_{SV}^2}$ Radial Anisotropy.

G, Ψ_G Azimuthal Anisotropy of V_{SV} , also related to SKS splitting (when horizontal symmetry axis).

+ a priori information (from mineralogy, ...)

- Body Waves (*Crampin, 1984*)

$$\rho V_{qSV}^2 = L + G_c \cos 2\Psi + G_s \sin 2\Psi$$

$$\rho V_{qSH}^2 = N - E_c \cos 4\Psi - E_s \sin 4\Psi$$

Geodynamic Interpretation

Convective cell: anisotropic parameters

S- Velocity

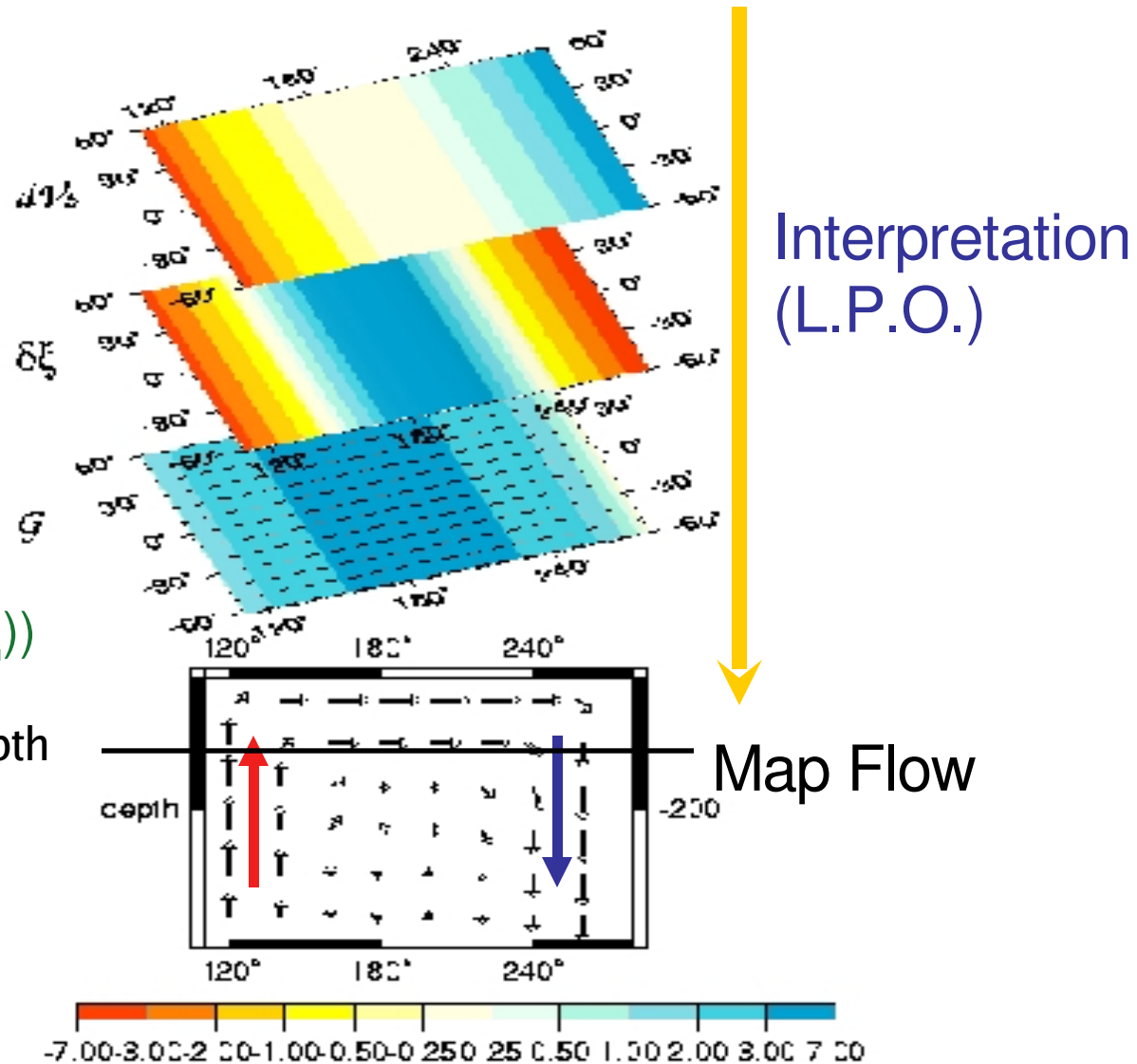
Radial Anisotropy

$$\xi = (V_{SH}^2 - V_{SV}^2) / V_{SV}^2$$

Azimuthal Anisotropy

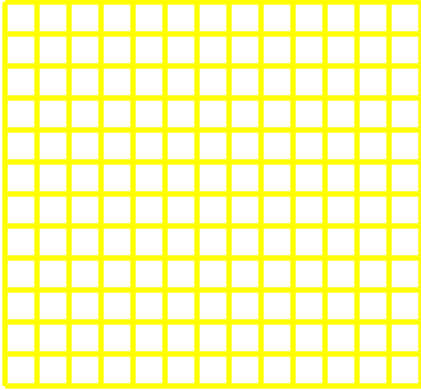
$$V_{SV} \approx V_{SV0} + \frac{1}{2} G \cos(2(\Psi - \Psi_G))$$

At a given depth



2 D tomography: N cells

Isotropic Inversion:



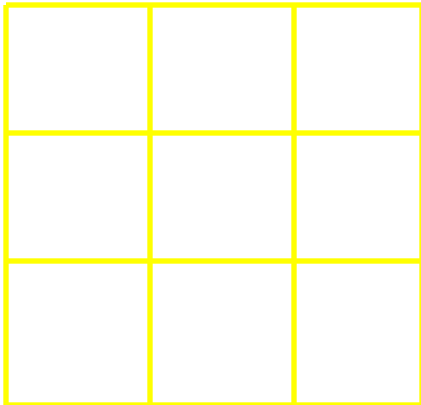
N independent parameters

0- Ψ term

VR1

Variance reduction

Anisotropic inversion



$$3N' = N$$

0+2 Ψ term

VR2

$VR2 > VR1 \Rightarrow$ the anisotropic model can be simpler than the isotropic model

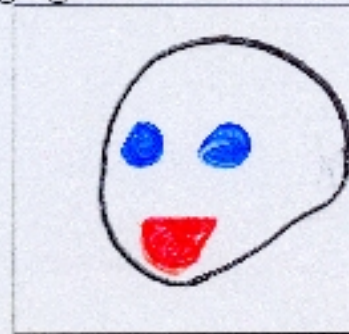
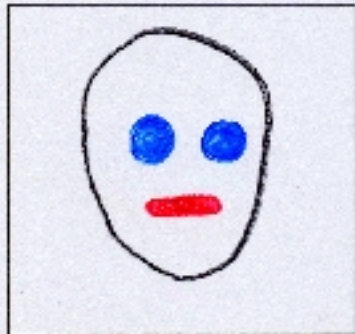
IMAGING OF FAMOUS SCIENTISTS

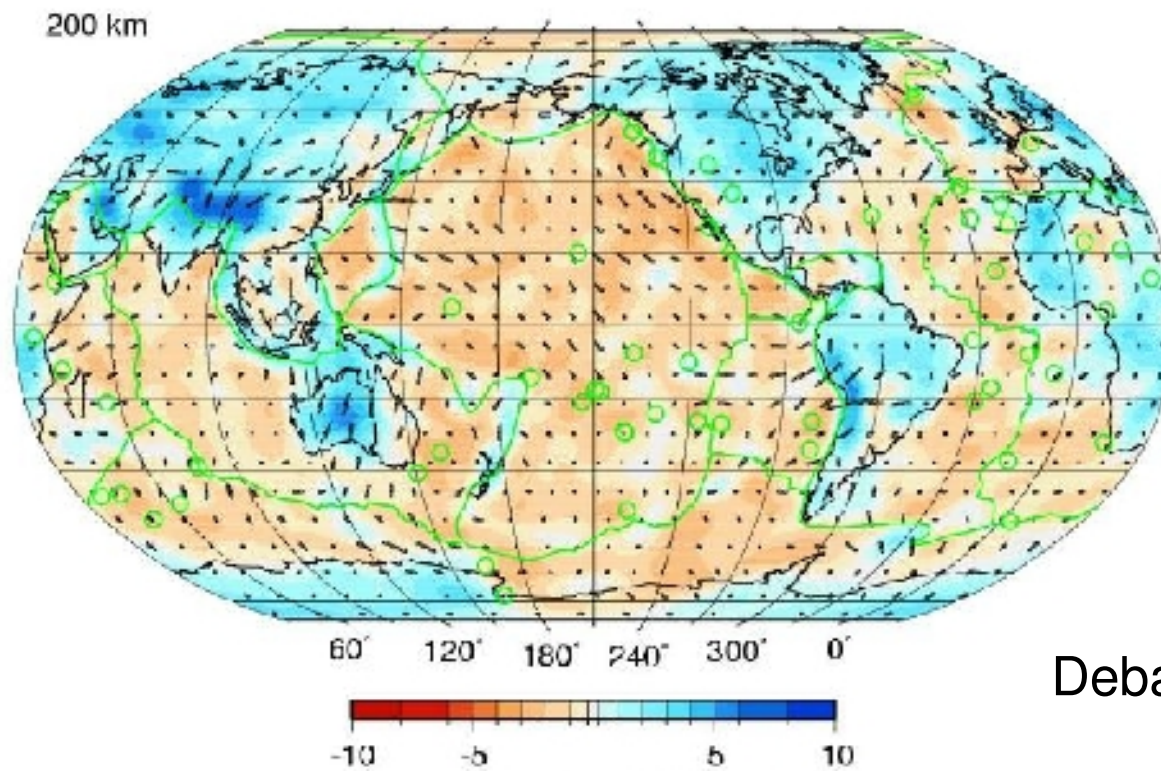
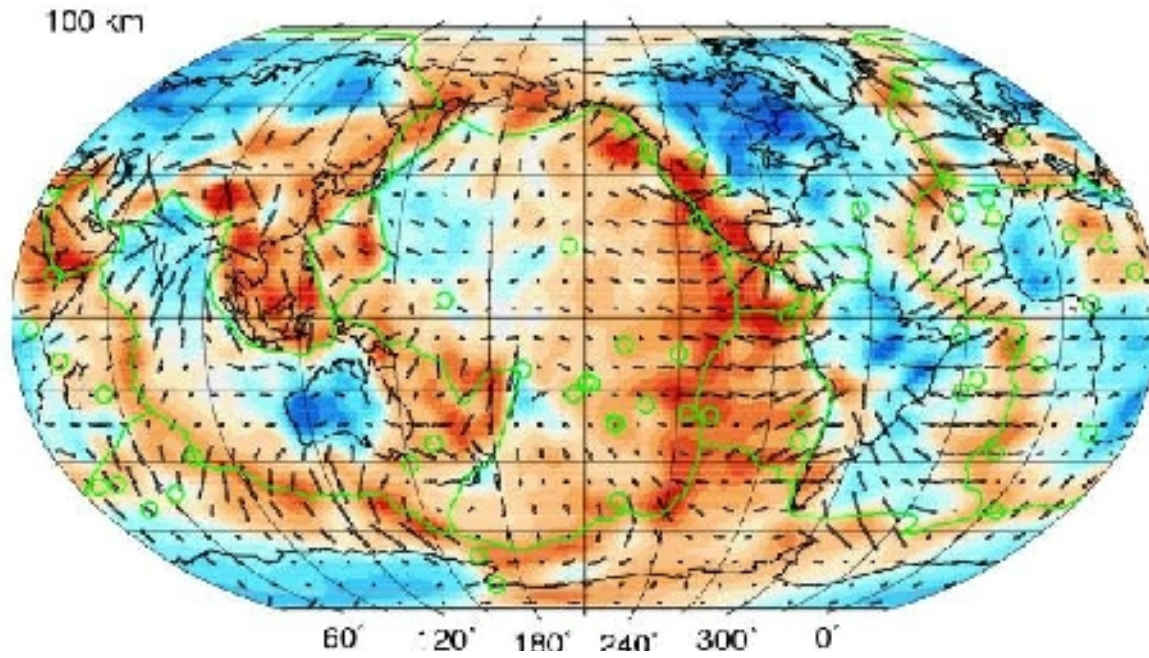


Anisotropic imaging



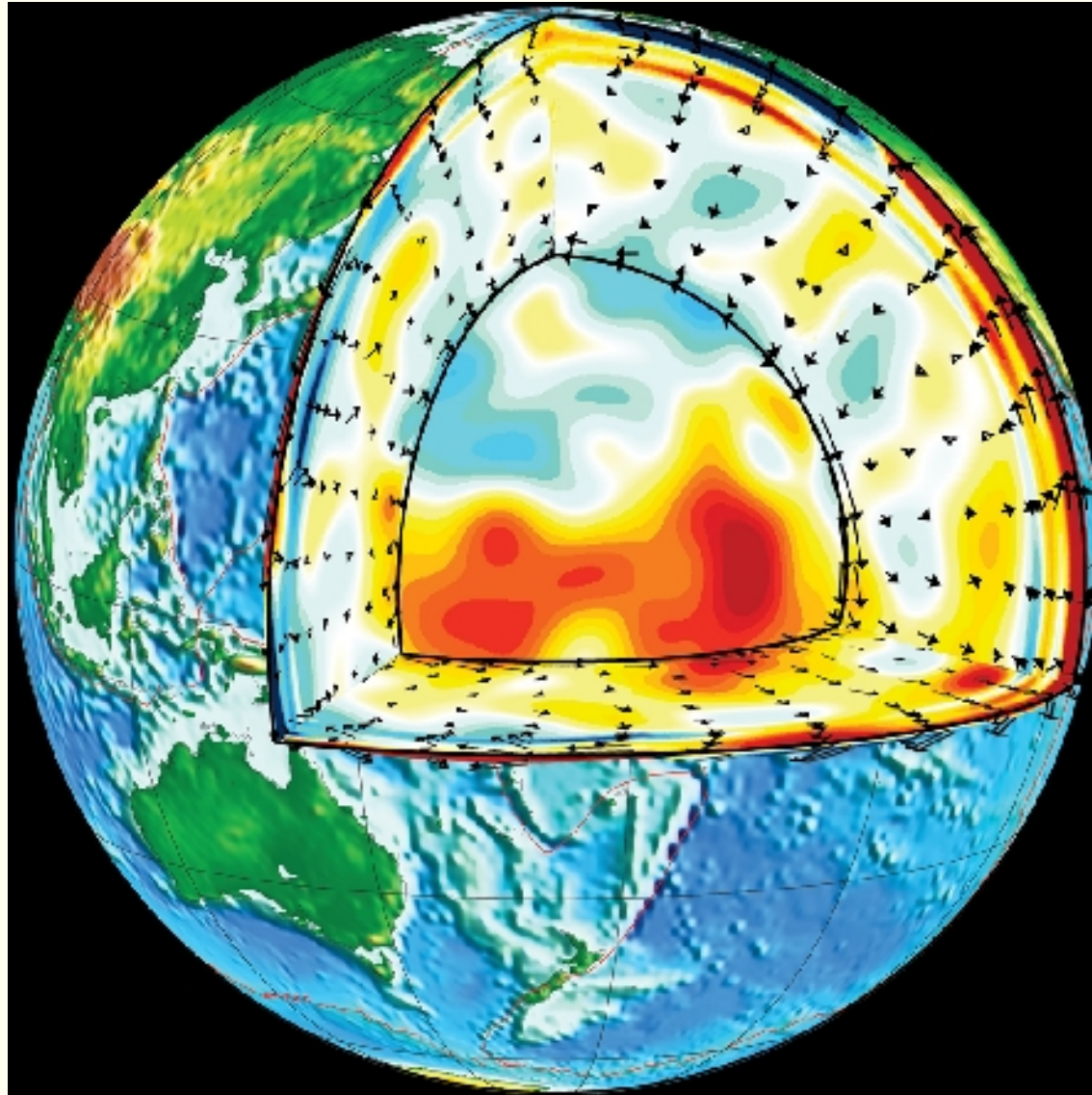
Isotropic Imaging





Debayle et al., 2005

Anisotropy- Geodynamics Relationship



CONCLUSIONS

- Progrès en instrumentation (Fond de Mer, Planète Mars; Spatial)
- Théorie des rais - Modes Normaux -> Méthodes numériques de plus en plus puissantes et précises grace a des ordinateurs de plus en plus puissants.
- Du Global vers le régional- Incorporation de nouveaux paramètres (anisotropie, anélasticité) en tomographie
- Approche pluridisciplinaire systématique: Confrontation des résultats sismologiques, avec expériences analogiques et numériques